Nadprzewodniki na bazie żelaza w świetle badań spektroskopią mössbauerowską

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Superconductivity in the non-magnetic state of iron under pressure

K. Shimizu et al. Nature 412, 316 (2001)

hcp Fe

becomes superconductor at temperatures below **2** K and at pressures between **15** and **30 GPa**



History of Superconductivity

introduction



Yoichi Kamihara,*,† Takumi Watanabe,‡ Masahiro Hirano,†,§ and Hideo Hosono†,‡,§



Layered Structure of Fe-based Superconductors



Mössbauer Spectroscopy



 γ -ray energy is modulated by the Doppler effect due to the source motion vs. absorber







Hyperfine Interactions between Nuclei and Electrons \Rightarrow Mössbauer Parameters



Electric Quadrupole Interaction \Rightarrow Quadrupole Splitting \Rightarrow **Electric Field Gradient EFG**



Magnetic Dipole Interaction \Rightarrow Magnetic Splitting \Rightarrow Magnetic hyperfine field *B*



Hyperfine Interactions



Magnetic Hyperfine Field

Electric Field Gradient + Magnetic Hyperfine Field





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Fe

Fe_{1.01}Se $T_{sc} = 8 \text{ K}$



Hyperfine magnetic field is equal to applied external magnetic field - it means that there is no magnetic moment on the Fe atoms

Mössbauer spectroscopy evidence for the lack of iron magnetic moment in superconducting FeSe

A.Błachowski, K.Ruebenbauer, J.Żukrowski, J.Przewoźnik, K.Wojciechowski, Z.M.Stadnik JALLCOM 494, 1 (2010)



Crystal structure of FeAs

- 1. Orthorhombic structure
- 2. The *Pnma* symmetry group
- 3. Arrows show $Pna2_1$ distortion

[0 *k*+1/2 0] iron and [0 *k* 0] iron

Magnetic structure of FeAs

magnetic moments lay in a-b plane, make antiferromagnetic spirals propagating along the c-axis and incommensurate with the lattice period





A.Błachowski, K.Ruebenbauer, J.Żukrowski, Z.Bukowski, Magnetic anisotropy and lattice dynamics in FeAs studied by Mössbauer spectroscopy, J. Alloys Compd. **582**, 167 (2014)

Anisotropy of the hyperfine magnetic fields (spiral projections onto *a-b* plane) in FeAs

Left column shows $\begin{bmatrix} 0 & k+1/2 & 0 \end{bmatrix}$ iron, right column shows $\begin{bmatrix} 0 & k & 0 \end{bmatrix}$ iron.

 B_a and B_b - iron hyperfine field components along the *a*-axis and *b*-axis, respectively.



"122" Fe-based Superconducting Family

<pre>AFe₂As₂ (A = Ca, Sr, Ba, Eu)</pre>				
Parent compound				
Ba	Fe ₂ As ₂ T _{SDW} = 136 K Superconductors			
Ba _{1-x} K _x Fe ₂ As ₂	Ba(Fe _{1-x} Co _x) ₂ As ₂	BaFe ₂ (As _{1-x} P _x) ₂		
<u>hole-doping</u> x = 0.40 T = 38 K	electron-doping x = 0.08 T = 24 K	<u>isovalent-substitution</u> x = 0.31 T = 31 K		
r_{SC} r	$F_{SC} = 2 + K$	Ba(Fe _{1-x} Co _x) ₂ As ₂		
	~	Rev. Mod. Phys. 87, 855 (2015)		

Charge density wave (CDW) - spatial modulation of the electron charge density

Electric field gradient wave (EFGW) - spatial modulation of the electric field gradient

Spin density wave (SDW) - spatial modulation of the electron spin density



The Mössbauer spectroscopy is sensitive to the spin and charge (electron) distribution around the resonant atoms via the hyperfine magnetic field, isomer shift and electric quadrupole interaction.

Spin density wave (SDW) as seen by Mössbauer Spectroscopy

$$B(qx) = \sum_{n=1}^{N} h_{2n-1} \sin[(2n-1)qx]$$

- h_{2n-1} amplitudes of subsequent harmonics
- q wave number of SDW
- *x* relative position of the resonant nucleus along propagation direction of SDW



CDW, EFGW and SDW as seen by Mössbauer Spectroscopy



For CDW one can estimate dispersion (around average value) of the electron density $\Delta \rho$

 $\Delta \rho = \sqrt{(\Gamma^2 - \Gamma_0^2) / \alpha^2} \qquad \qquad \Gamma_0 = 0.1 \text{ mm/s} \qquad - \text{ unbroadened line width} \\ \alpha = -0.291 \text{ (mm/s) (Bohr)}^3 \text{ el.}^{-1} \qquad - \text{ calibration constant}$



⁵⁷Fe Mössbauer spectra





 $\pi/2$





⁵⁷Fe Mössbauer spectra



Compound	CaFe ₂ As ₂	BaFe ₂ As ₂	EuFe ₂ As ₂
T_c (K)	175.3(3)	136.0(1)	192.1(1)
α ₀	0.158(2)	0.102(1)	0.124(1)

critical exponent $\alpha_0 \approx 0.125 \implies$ universality class (1, 2)

one dimension in the spin space (Ising model) and

two dimensions in the real space (magnetic planes)

A.Błachowski, K.Ruebenbauer, J.Żukrowski, K.Rogacki, Z.Bukowski, J.Karpinski, Shape of spin density wave versus temperature in AFe₂As₂ (A = Ca, Ba, Eu) Phys. Rev. B **83**, 134410 (2011)



⁵⁷Fe Mössbauer spectra



SDW is suppressed by doping

A.Błachowski, K.Ruebenbauer, J.Żukrowski, Z.Bukowski, M.Matusiak, J.Karpinski, Interplay between spin density wave and superconductivity in '122' iron pnictides Acta Phys. Pol. A **121**, 726 (2012)





A.Błachowski, K.Ruebenbauer, J.Żukrowski, Z.Bukowski, K.Rogacki, P.J.W.Moll, J.Karpinski, Interplay between magnetism and superconductivity in EuFe_{2-x}Co_xAs₂ studied by ⁵⁷Fe and ¹⁵¹Eu Mössbauer spectroscopy, Phys. Rev. B **84**, 174503 (2011)

¹⁵¹Eu Mössbauer spectroscopy



Eu(2+)

EuFe_{2-x}Co_xAs₂



EuFe_{2-x}Co_xAs₂

⁵⁷Fe Mössbauer spectra





Ba_{1-x}K_xFe₂As₂



A.K.Jasek, K.Komędera, A.Błachowski, K.Ruebenbauer, Z.Bukowski, J.G.Storey, J.Karpinski, Electric field gradient wave (EFGW) in iron-based superconductor Ba_{0.6}K_{0.4}Fe₂As₂ studied by Mössbauer spectroscopy, J. Alloys Compd. **609**, 150 (2014) ⁵⁷Fe Mössbauer spectra of BaFe₂As₂ parent compound and Ba_{0.6}K_{0.4}Fe₂As₂ superconductor



Inset shows electronic specific heat coefficient of superconductor.

⁵⁷Fe Mössbauer spectra of the $Ba_{0.6}K_{0.4}Fe_2As_2$ ($T_{sc} = 38$ K) across transition to the superconducting state.



$Ba_{0.6}K_{0.4}Fe_2As_2$ ($T_{SC} = 38$ K)

Mössbauer parameters:

- S spectrum shift versus α -Fe
- Δ_0 constant component of quadrupole splitting
- Γ absorber line width

Shape of **EFGW** *electric field gradient wave* (**d** electrons density variation)





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Magdalena Piskorz WFiIS AGH, Fizyka Techniczna, III rok

Czułość spektroskopii mössbauerowskiej na przejścia nadprzewodzące w pniktydkach żelaza 51 Konferencja Studenckich Kół Naukowych AGH maj 2014 WYRÓŻNIENIE

SmFeAsO_{1-x}F_x doping superconductivity parent compound $T_{sc} = 55K$ **SmFeAsO** F • T. 200 140 T_{mag} o ΔT_s Tetra. Δ.....Δ 120 ▼ T_{Sm} PM Temperature (K) 100 emperature (K) Tetra. 80 100 PM Fe 77 60 As Ortho. AFFe Static 40 magnetism Superconductivity SC 20 AFsm 0 0 0.00 0.06 80.0 0.04 0.02 0 0.10 0.20 0.30 x F doping

Y. Kamihara *et al.*, New J. Phys. **12**, 033005 (2010)

A.J. Drew et al., Nature Materials 8, 310 (2009)

our sample – SmFeAsO_{0.91}F_{0.09}

$SmFeAsO_{0.91}F_{0.09}$



-2

-1

0

Velocity (mm/s)

2

A.K.Jasek, K.Komędera, A.Błachowski, K.Ruebenbauer, H.Lochmajer, N.D.Zhigadlo, and K.Rogacki, J. Alloys Compd. **658**, 520 (2016)

⁵⁷Fe Mössbauer spectra of SmFeAsO_{0.91} $F_{0.09}$ ($T_{sc} \approx 47$ K) across transition to the superconducting state





Comparison between

charge density modulation changes during superconducting transition in



Comparison between

charge density modulation changes during superconducting transition in





Nematic phase is characterized by electronic anisotropy in *a-b* plane with broken rotational symmetry but preserved translational symmetry (tetragonal phase).







Conclusions

Precursor compound: FeAs

Antiferromagnetic incommensurate spirals propagate along c-axis with the hyperfine field varying along the spiral in a fashion resembling symmetry of 3d electrons (local magnetic moment) in a-b plane with significant distortion caused by arsenic bonding p electrons.

Parent compounds: AFe₂As₂

Longitudinal spin density wave SDW develops with triangular shape along propagation direction, i.e. a-axis. Upon temperature lowering they transform into rectangular form. Hence, they are indistinguishable from simple anti-ferromagnetic ordering close to the ground state.

<u>Superconductors:</u> $Ba_{1-x}K_xFe_2As_2$ and $SmFeAsO_{1-x}F_x$

The Mössbauer spectroscopy is sensitive to the superconducting transition in Fe-based superconductors via **<u>change of the electron charge density modulation</u>**, which is seen via dispersion of isomer shift (CDW caused by s electrons) and via distribution of electric field gradient (EFGW caused by d electrons).

Shape and amplitude of EFGW and CDW are strongly perturbed at the superconducting transition. Namely, all modulations are strongly changed at critical temperature due to the superconducting gap opening and subsequent formation of Cooper pairs. However dispersion of the charge density and EFGW shape behave in the opposite ways for these two superconductors.

Thank you very much for your attention !