Nowe pomiary momentów magnetycznych naładowanych leptonów

Mateusz Dyndał (KOiDC WFilS AGH)

Seminarium WFiIS AGH, 26.01.2024

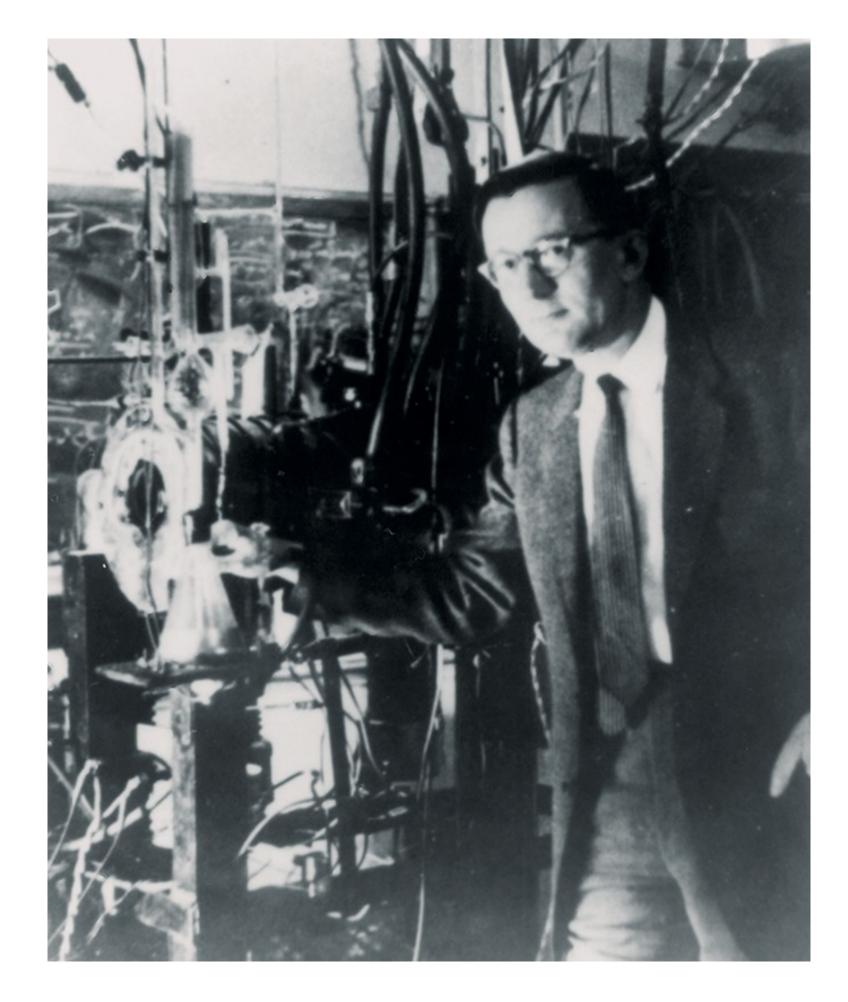






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Historical perspective

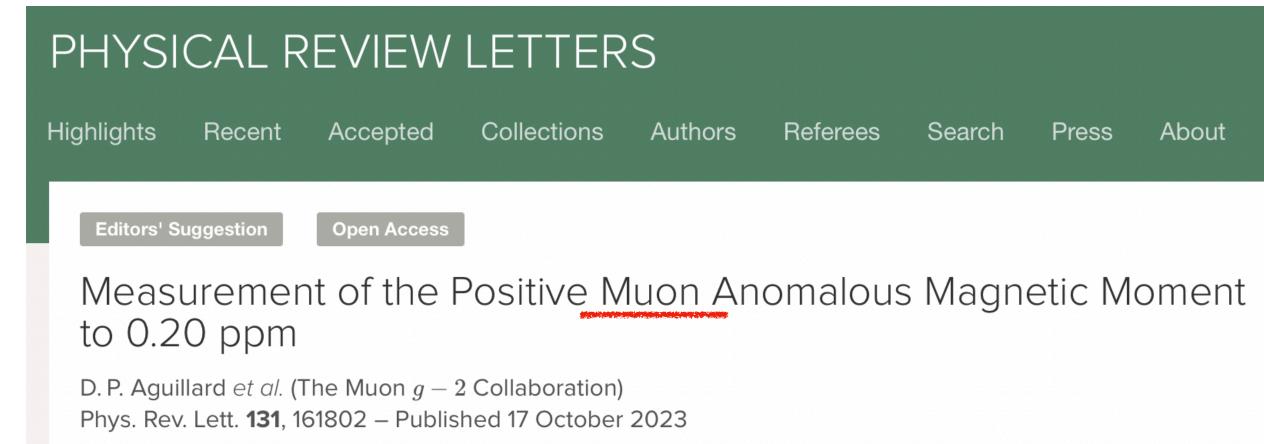


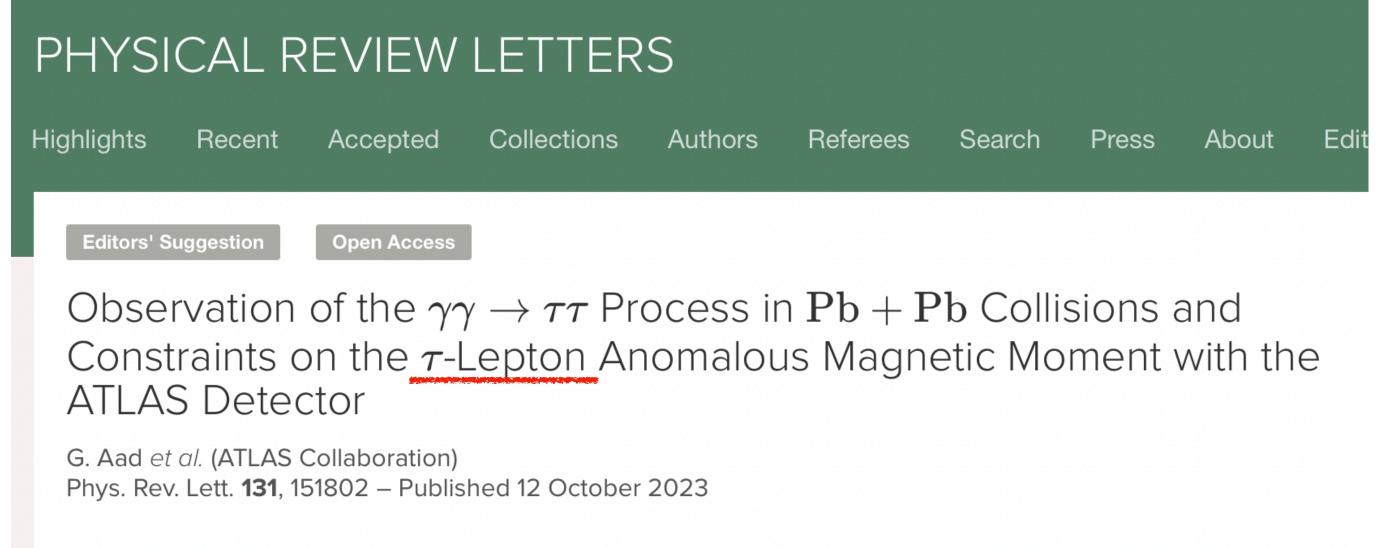
P. Kusch and H. M. Foley (1947)



New papers published in 2023







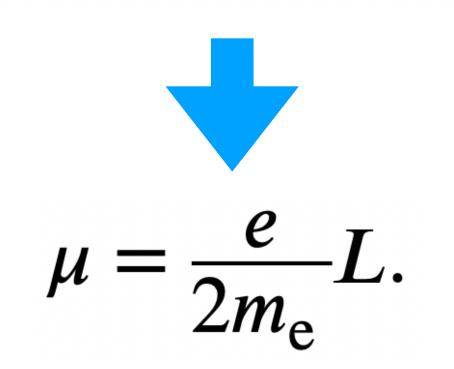
Magnetic (dipole) moment in classical physics

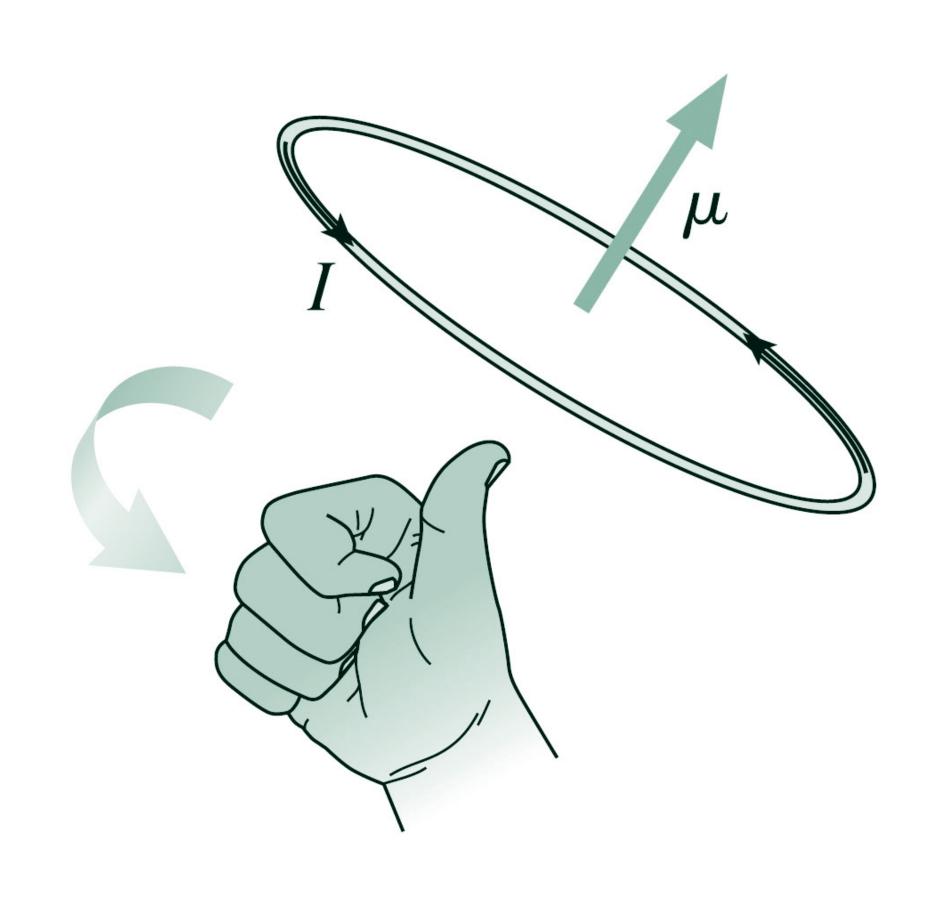
$$ec{M}=ec{\mu} imes ec{B}.$$

For single electron:

$$\mu = IA = \frac{e}{2\pi r / v} \pi r^2 = \frac{evr}{2}.$$

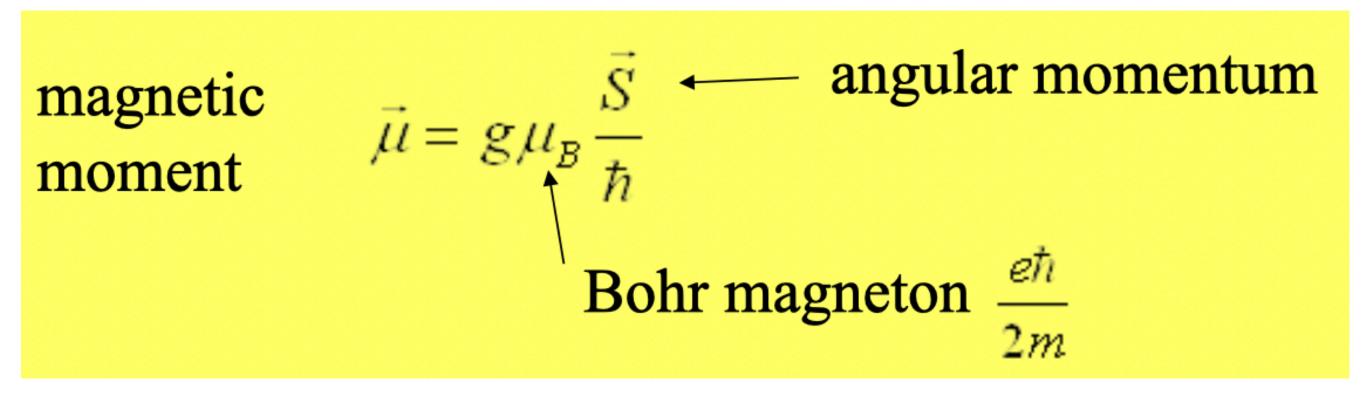
$$|\vec{L}| = |\vec{r}| \times |\vec{p}| = rp \sin \theta = rp = rm_e v.$$

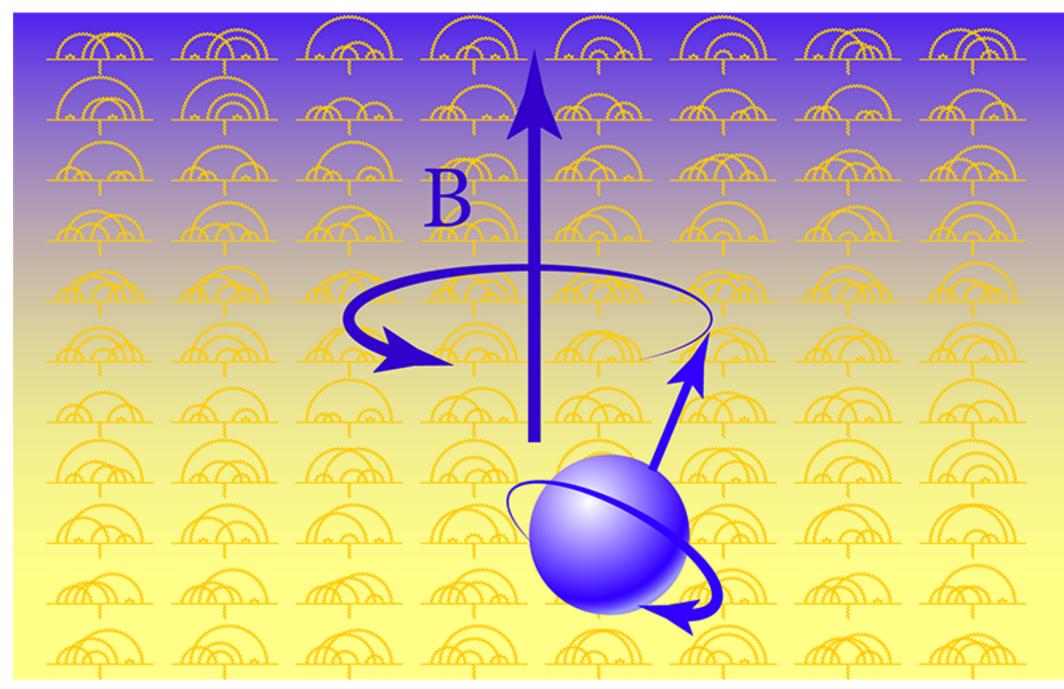




QM: Magnetic (dipole) moment and the g-factor

• g = proportionality constant between spin and magnetic moment

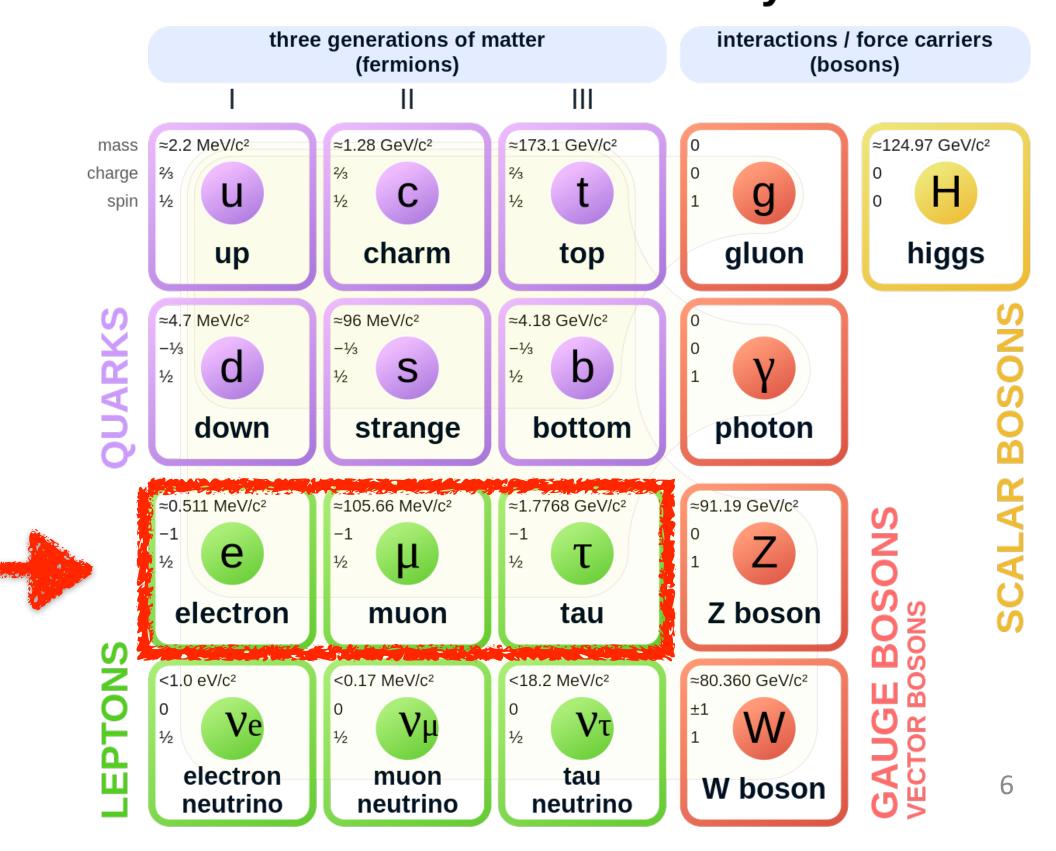




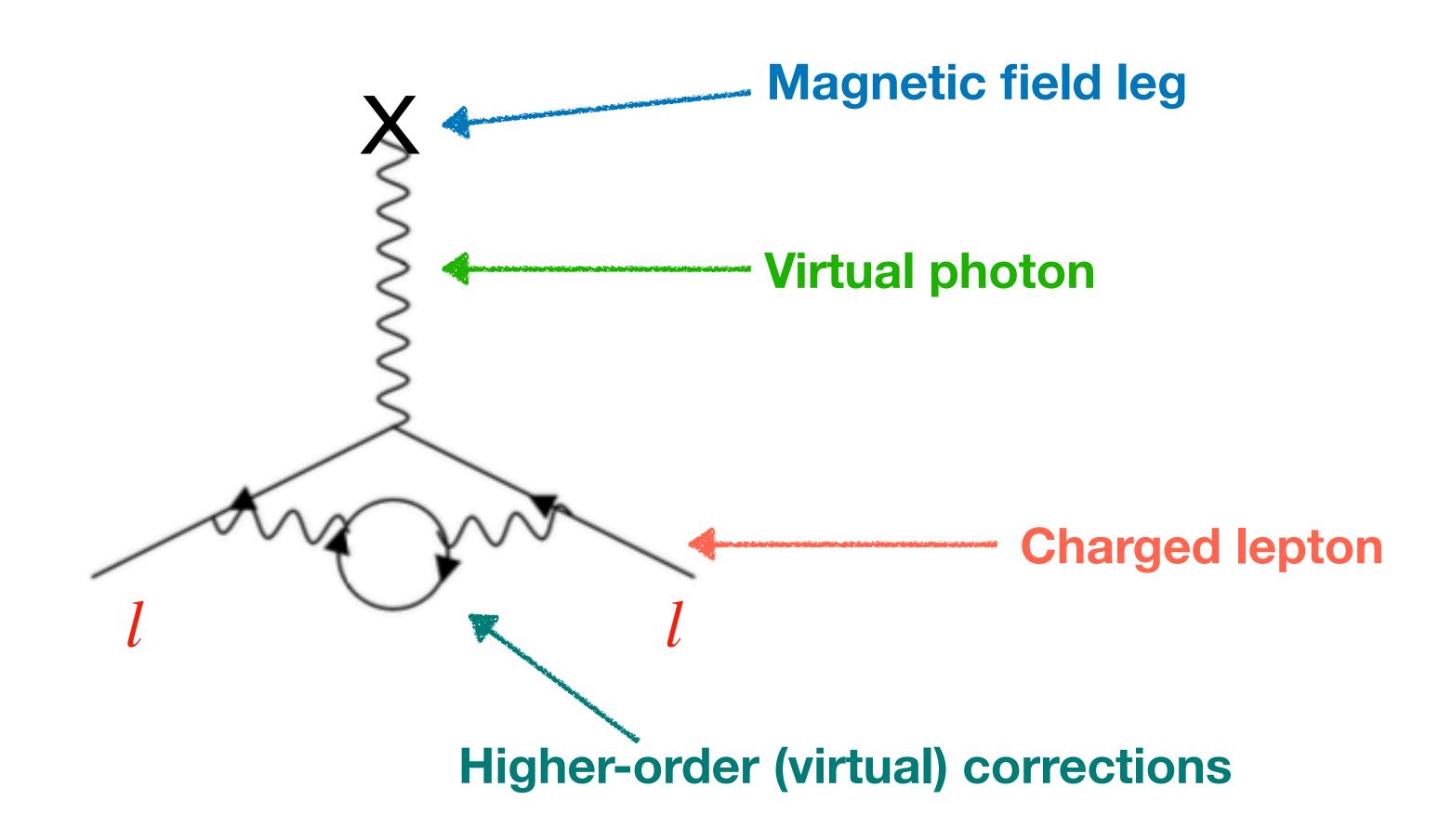
Charged leptons in the Standard Model

- Standard Model (SM) = our current understanding of fundamental interactions
- Predicts three generations of matter particles (fermions)
 - Including three types of charged leptons:
 - -> electron
 - -> muon
 - -> tau
 - Charged leptons interact ~identically in SM;
 they only differ in mass

Standard Model of Elementary Particles



Interaction of charged leptons with magnetic field in QED



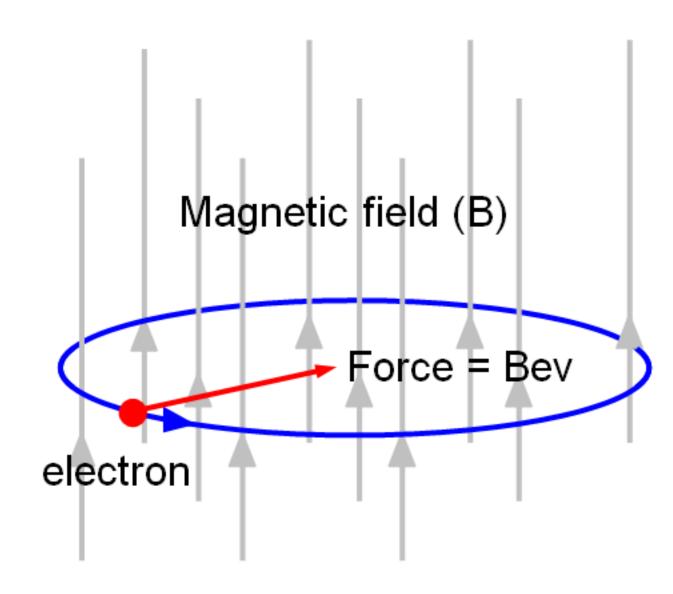
Magnetic (dipole) moment of charged leptons

- For electrons/muons/taus Standard Model predicts g ≈ 2
- Note: quantum fluctuations make it slightly larger than 2...
- Anomalous magnetic moment
 - $a_{\ell} = (g_{\ell} 2)/2$
 - Shows how much g differs fractionally from 2!

 a_{ℓ}

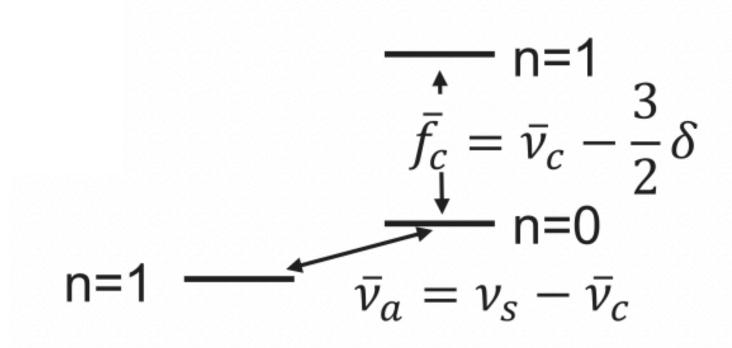
Electron in a magnetic field

Classically



$$\frac{mv^2}{r} = qBv \rightarrow \nu_c = eB/(2\pi m)$$

In QM



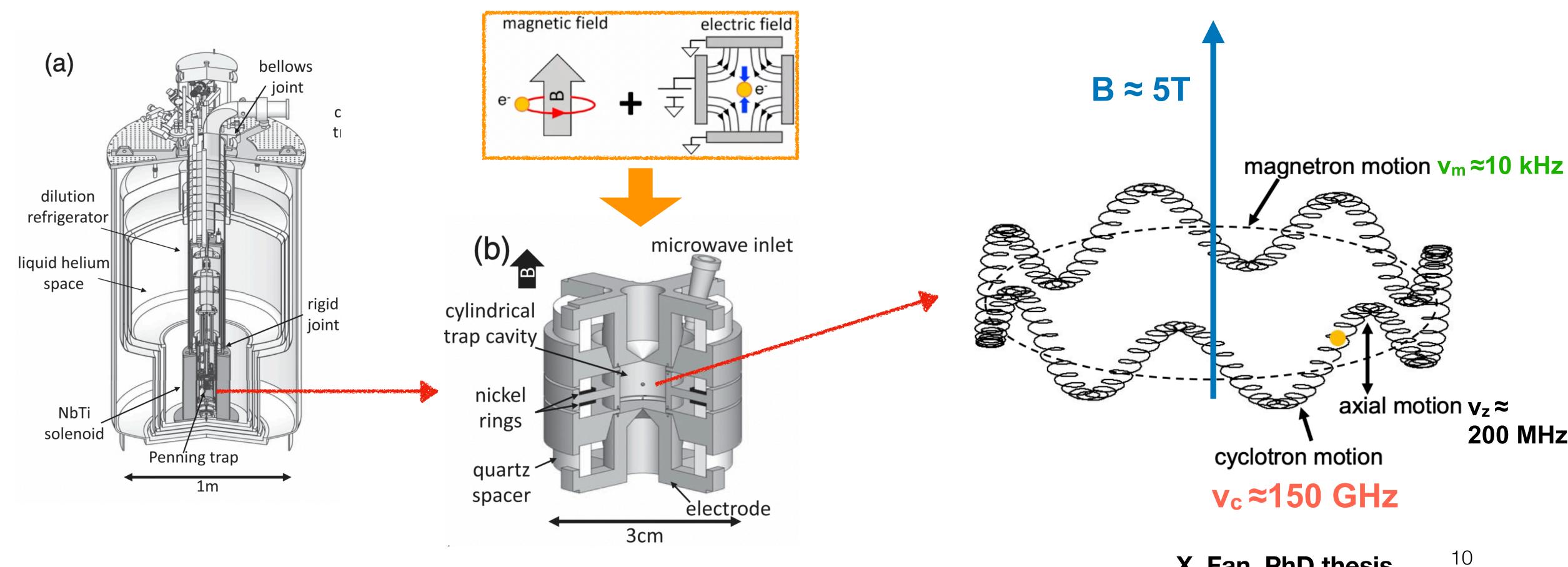
$$m_s = -1/2$$
 $m_s = +1/2$

$$E = h\nu_s m_s + h\nu_c \left(n + \frac{1}{2} \right)$$

$$\frac{g}{2} = \frac{\nu_s}{\nu_c} = 1 + \frac{\nu_a}{\nu_c}$$

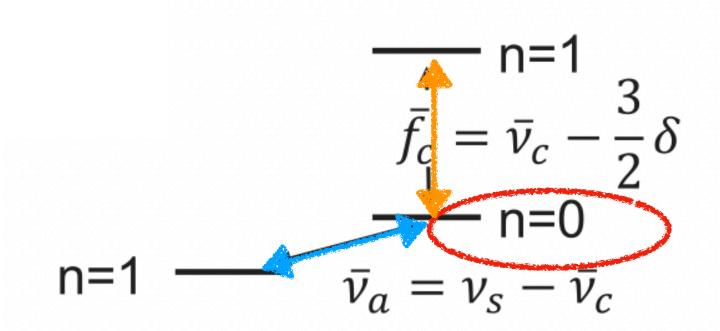
Electron g-2: measurement idea

Using one-electron quantum cyclotron (Penning trap), with T<100 mK



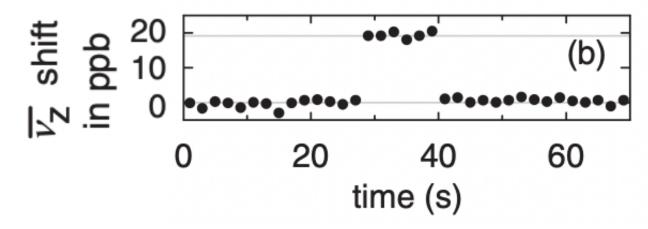
Electron g-2: Measurement procedure

- Cyclotron quantum jump spectroscopy
- Each cyclotron and anomaly quantum jump trial starts with preparing the electron in the spin-up ground state, |n=0, m_S=+1/2>

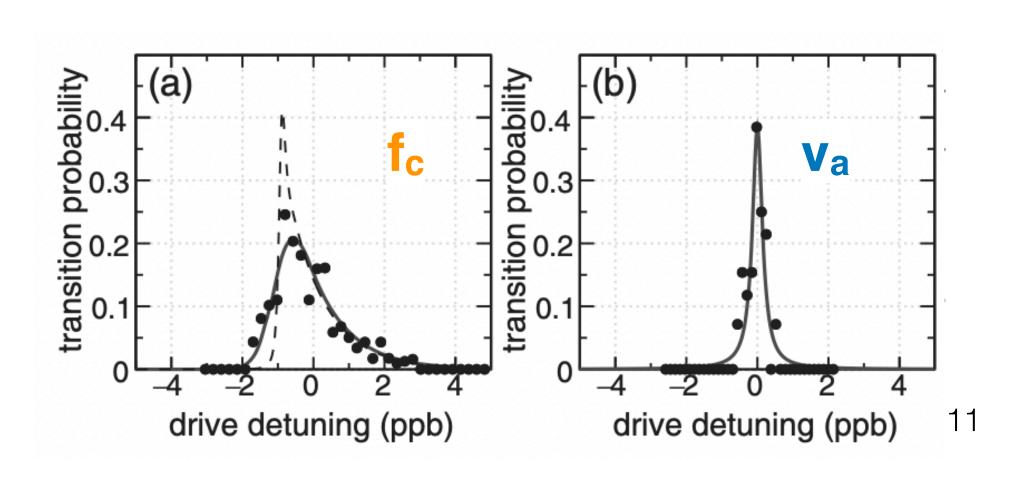


$$m_s = -1/2$$
 $m_s = +1/2$

Quantum jumps detected via shifts in the axial frequency



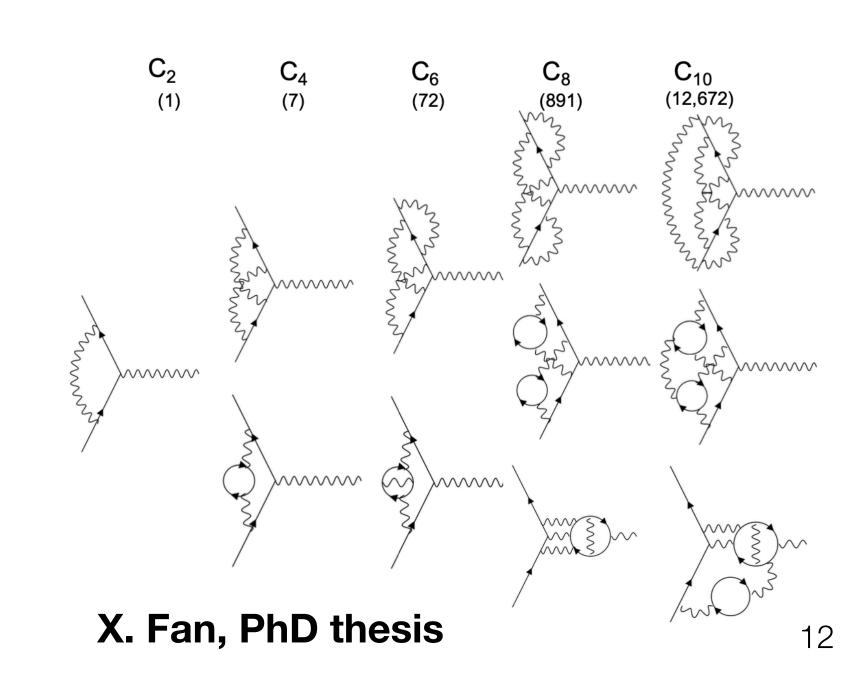
 Multiple attempts at different frequencies are binned in a histograms to reveal the frequency lines:



Electron g-2: status of theory calculations

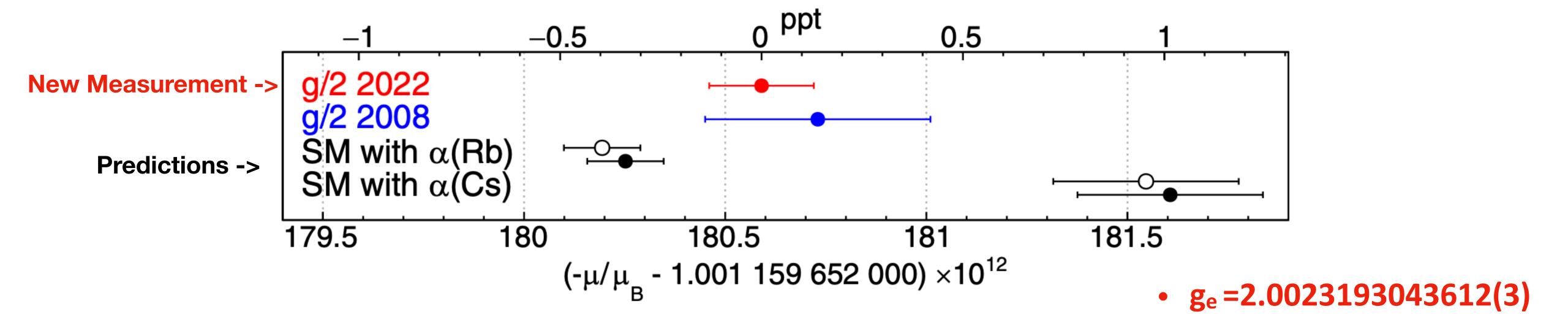
- QED provides the asymptotic series in powers of α , along with the muon and tau contributions ($a_{\mu\tau}$)
- The constants C₂, C₄, C₆, C₈ calculated exactly, but require measured lepton mass ratios as input
- The measurement is so precise that a numerically calculated tenth order C₁₀ is required
- Alternative theory evaluation of C₁₀
 differs slightly for reasons not yet understood

$$\frac{g}{2} = 1 + C_2 \left(\frac{\alpha}{\pi}\right) + C_4 \left(\frac{\alpha}{\pi}\right)^2 + C_6 \left(\frac{\alpha}{\pi}\right)^3 + C_8 \left(\frac{\alpha}{\pi}\right)^4 + C_{10} \left(\frac{\alpha}{\pi}\right)^5 + \dots + a_{\mu\tau} + a_{\text{hadronic}} + a_{\text{weak}}.$$



Electron g-2: Results

Phys.Rev.Lett. 130 (2023) 7, 071801



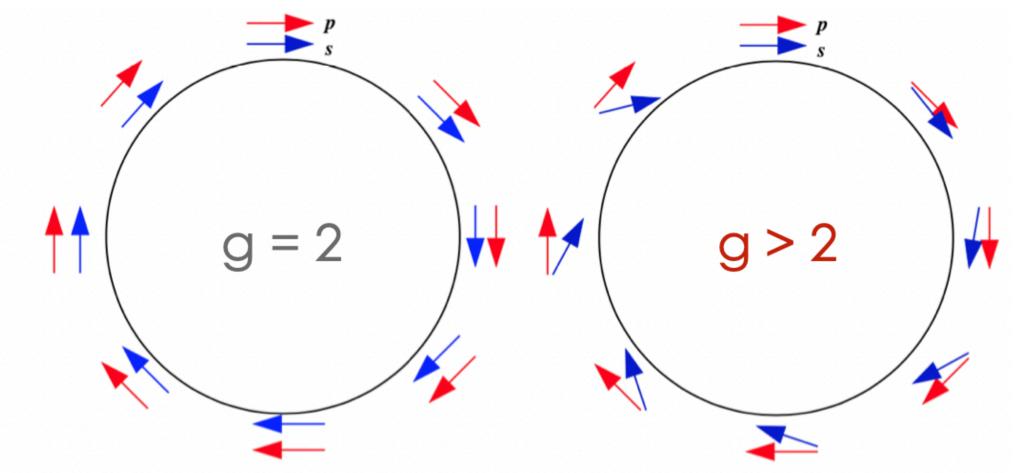
- Significant uncertainty reduction wrt previous measurement (blue), reaching **0.13 ppt** relative precision (ppt is 10⁻¹²)
- SM predictions (solid and open black points for slightly differing C₁₀ calculations) are functions of discrepant α measurements (independent measurements use Rb or Cs atom recoil)
- New measurement with positron (CPT symmetry test) underway

How to measure muon g-2?

- Why muon g-2?
 - The muon is 200 times more massive than electron \rightarrow (m_e/m_μ)² ~40 000 times more sensitive to new massive particles
- In a magnetic storage ring, the muon spin precesses slightly faster than the cyclotron frequency:

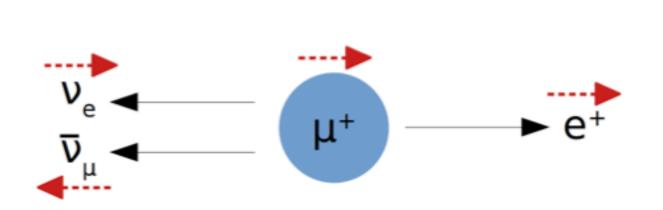
$$\underline{\vec{\omega}_s} = -\frac{ge\vec{B}}{2m} - (1 - \gamma)\frac{e\vec{B}}{m\gamma} \qquad \underline{\vec{\omega}_c} = -\frac{e\vec{B}}{m\gamma}$$

$$\vec{\omega}_a = \underline{\vec{\omega}_s} - \underline{\vec{\omega}_c} = -\left(\frac{g-2}{2}\right)\frac{e\vec{B}}{m} \equiv -\underline{a_\mu}\frac{e\vec{B}}{m}$$

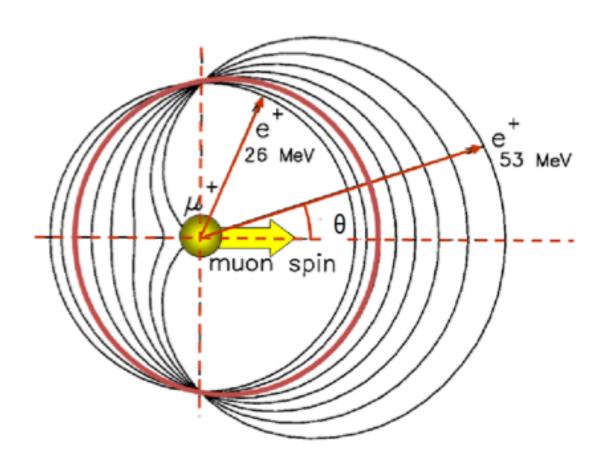


• Measure ω_a and $B \rightarrow$ obtain a_μ

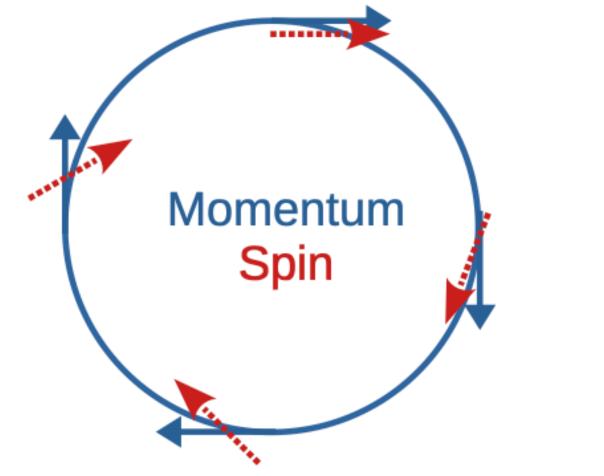
How to measure ω_a for muons?



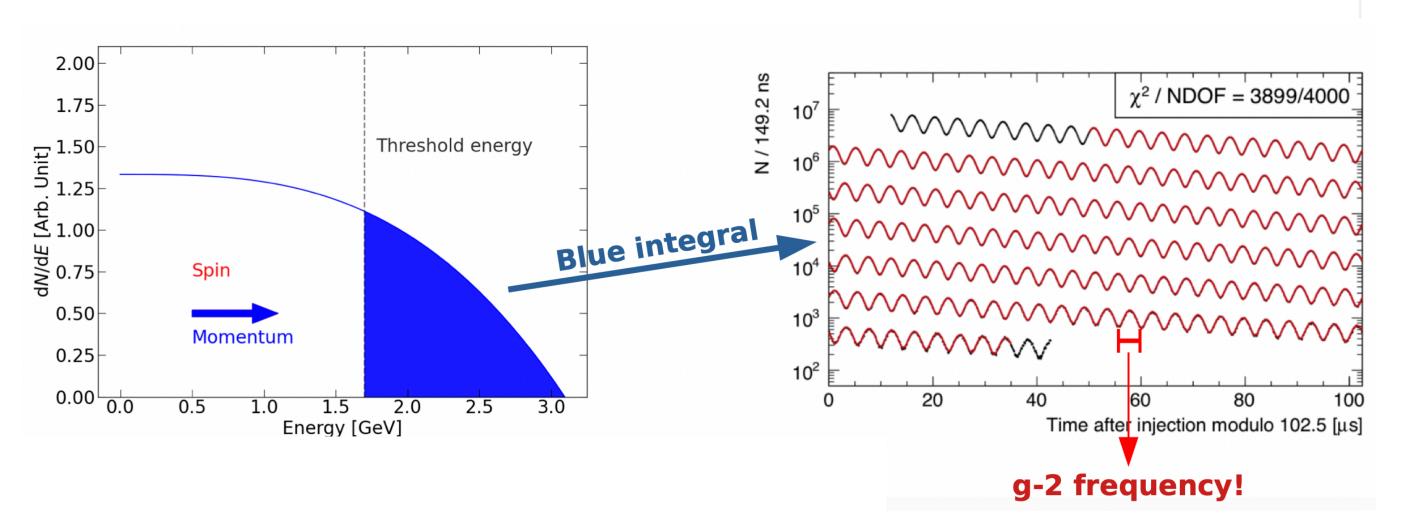
Muon decays in a positron and 2 neutrinos



Parity violation → positrons in CM preferably in the direction of the muon spin

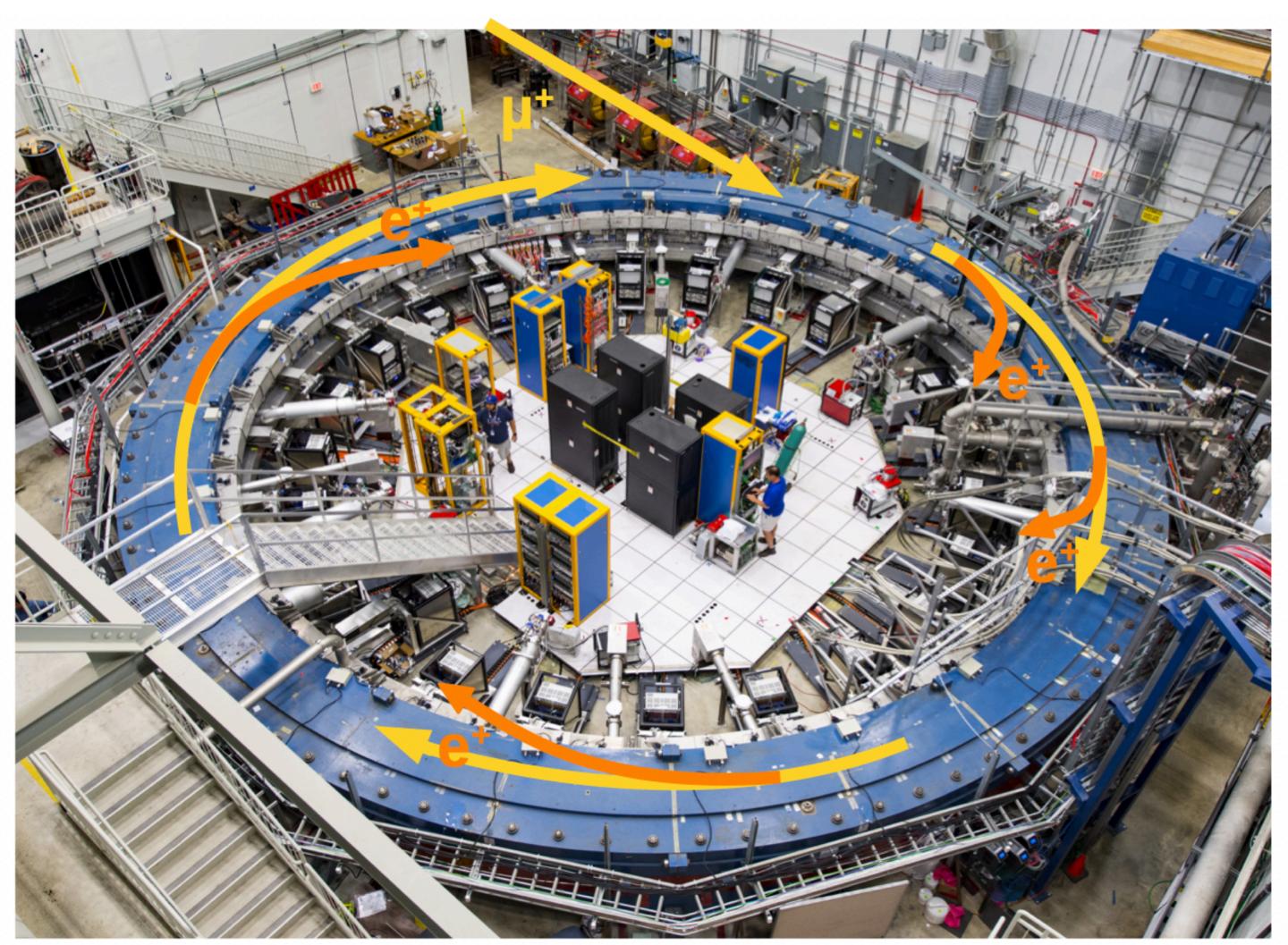


Spin precession → the energy spectrum in the lab frame oscillates through time

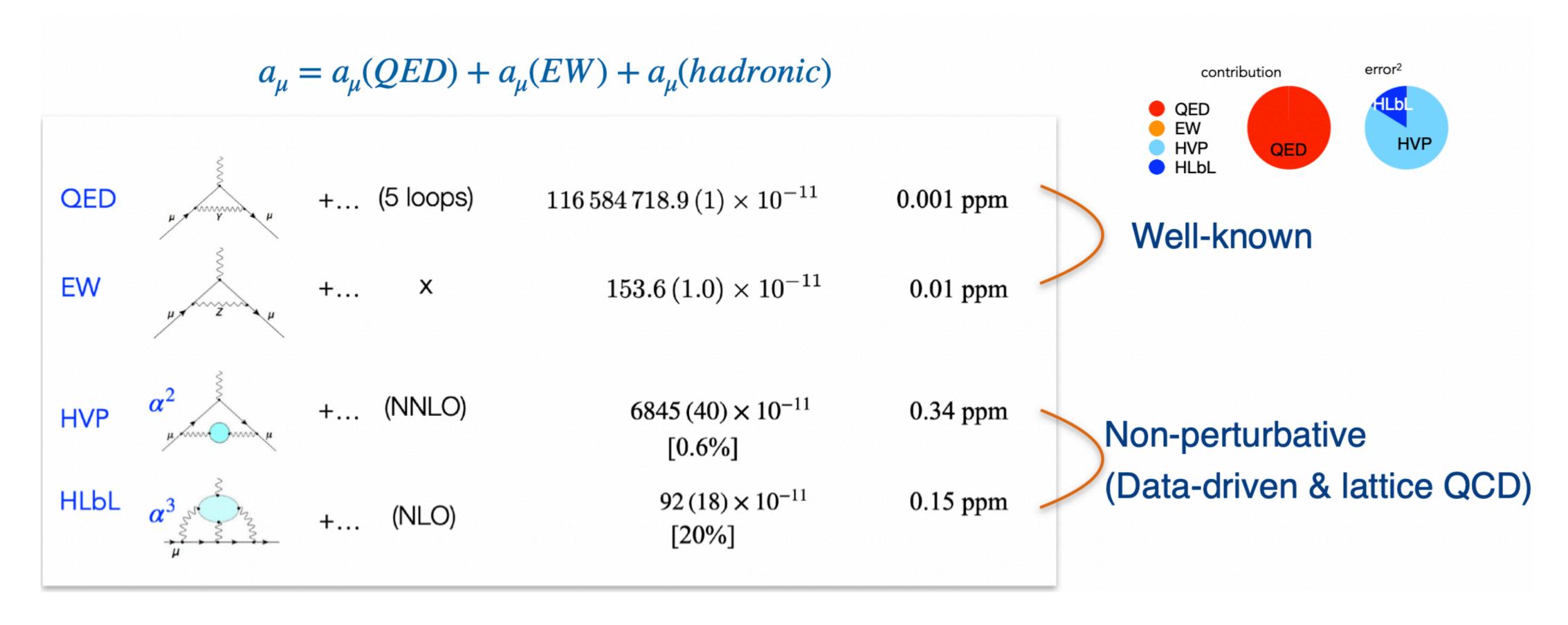


The Muon g-2 Experiment

- Located in Fermilab, Chicago (continuing the experiment conducted at Brookhaven)
- 15 m-diameter superconducting magnet with an exceptionally uniform magnetic field, used as a storage ring
- 24 EM calorimeters to measure decay positrons (on the inside of the storage ring)
- B-field value actively mapped using an NMR probe



Muon g-2: status of theory calculations



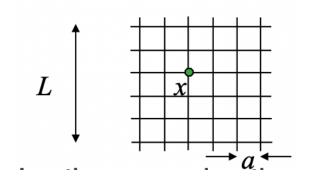
- QED and EW contributions are very well-known with small uncertainties
- Hadronic vacuum polarisation (HVP) contribution error dominates the uncertainty budget

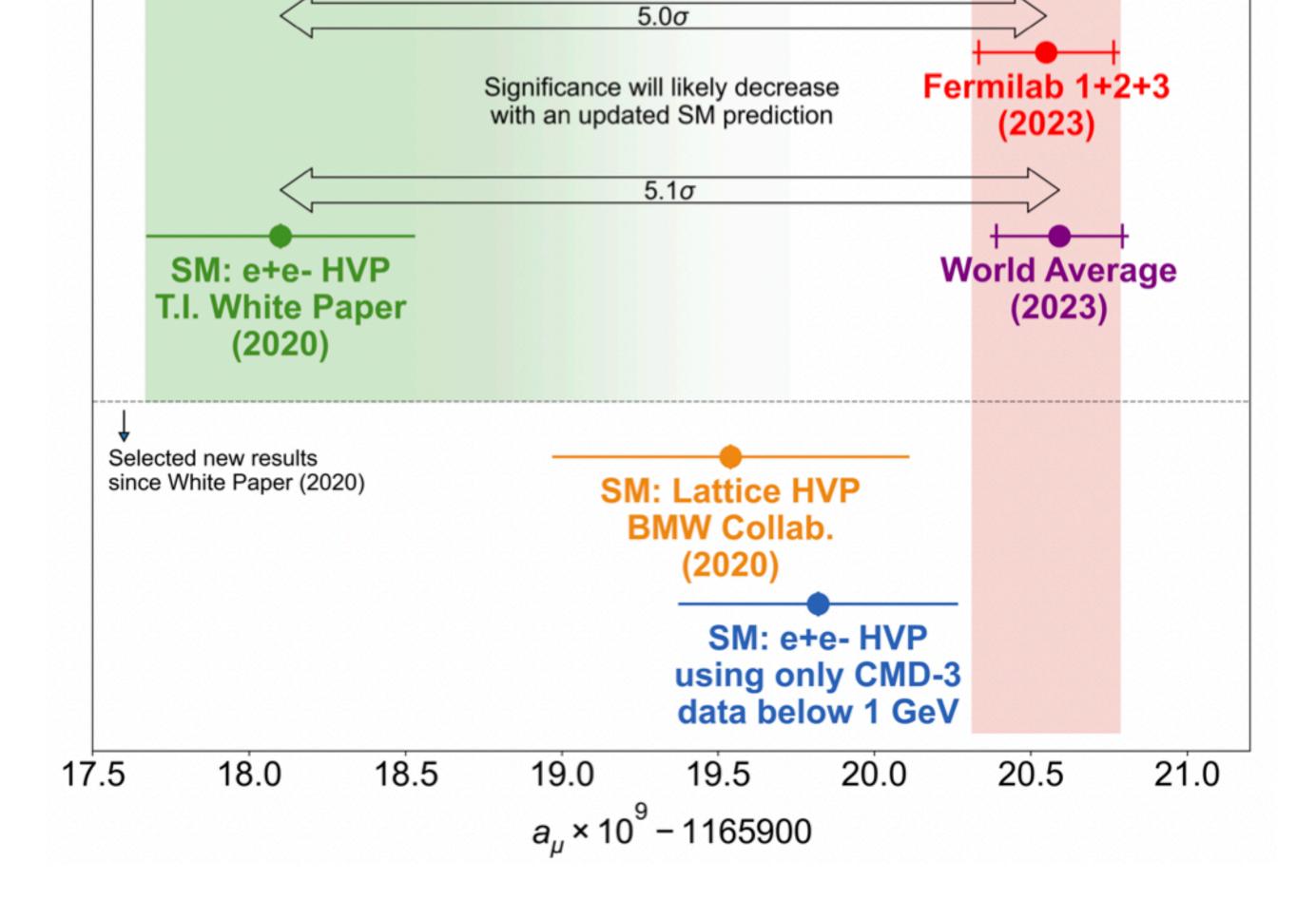
Muon g-2: most recent result

Phys.Rev.Lett. 131 (2023) 16, 161802

- Theory calculation differ in HVP calculations:
 - lattice QCD (Ab-initio)
 - dispersive (e+e-) method (data-driven)
 - this results in "theory-theory" tensions...

lattice

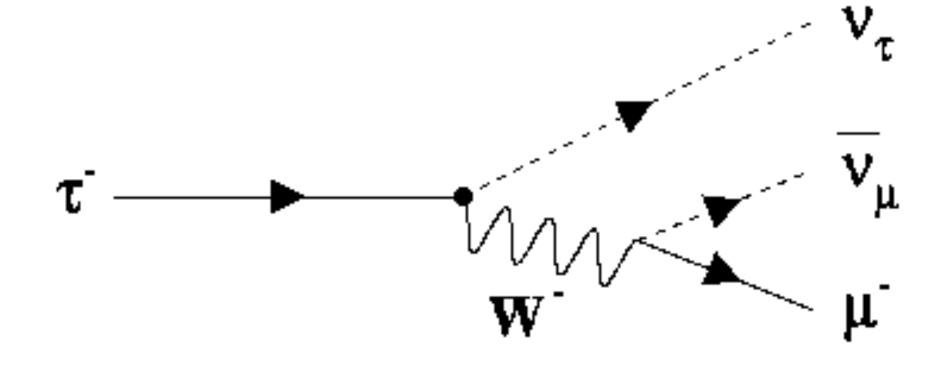






The tau lepton

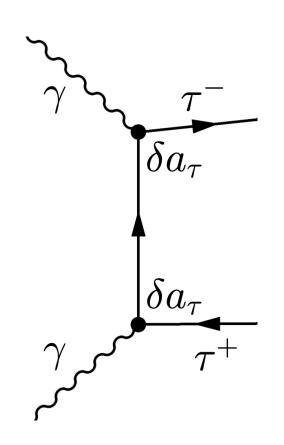
- Discovered in 1970's, it's the heaviest charged lepton
 - ≈ 2000 heavier than the electron
 - Due to large mass, it decays almost immediately (lifetime of 3×10-13 s)
- Because of extremely short lifetime impossible to make experiments with magnetic field interaction

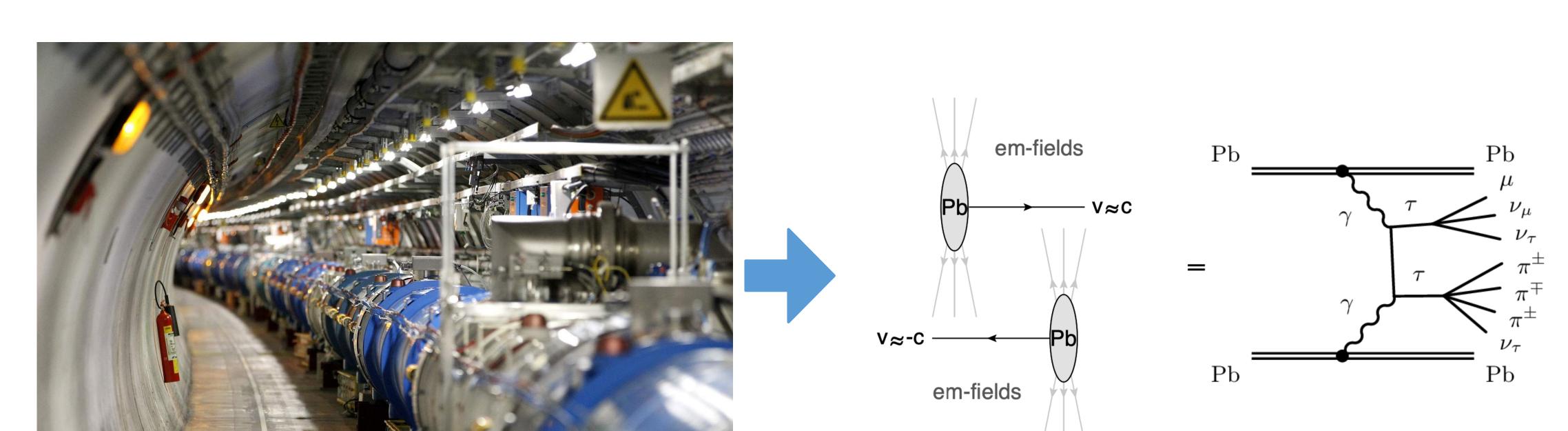


Example tau decay into neutrinos and a muon

Tau lepton EM interactions at the LHC

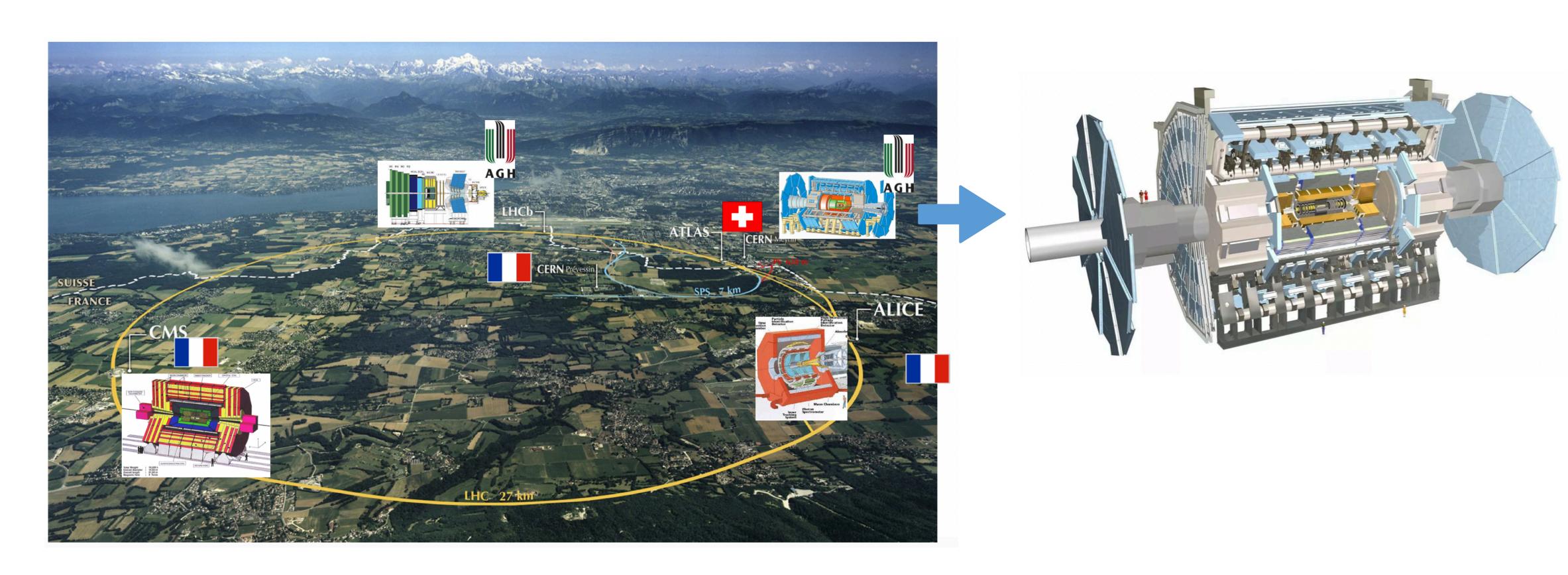
- We can measure the strength of EM interaction with tau lepton by studying the following process:
 - But: a powerful source of high-energy photons is needed...
- Heavy (charged) ions are intense source of photons
 - Fortunately we collide lead ions at the Large Hadron Collider!





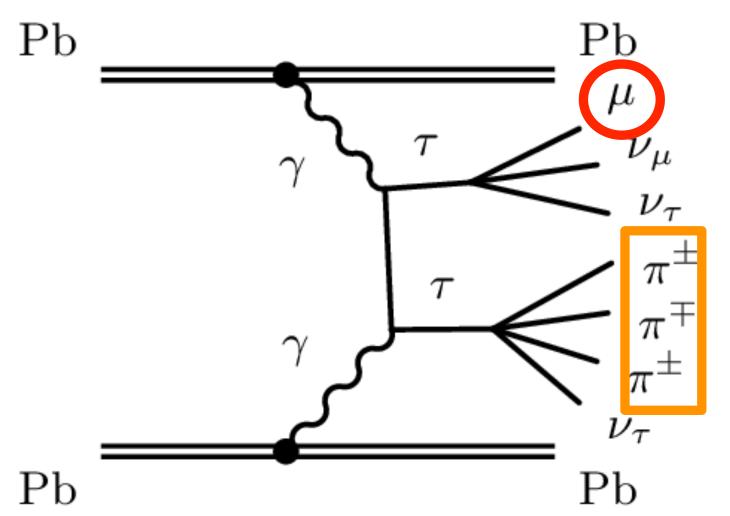
Large Hadron Collider and the ATLAS experiment

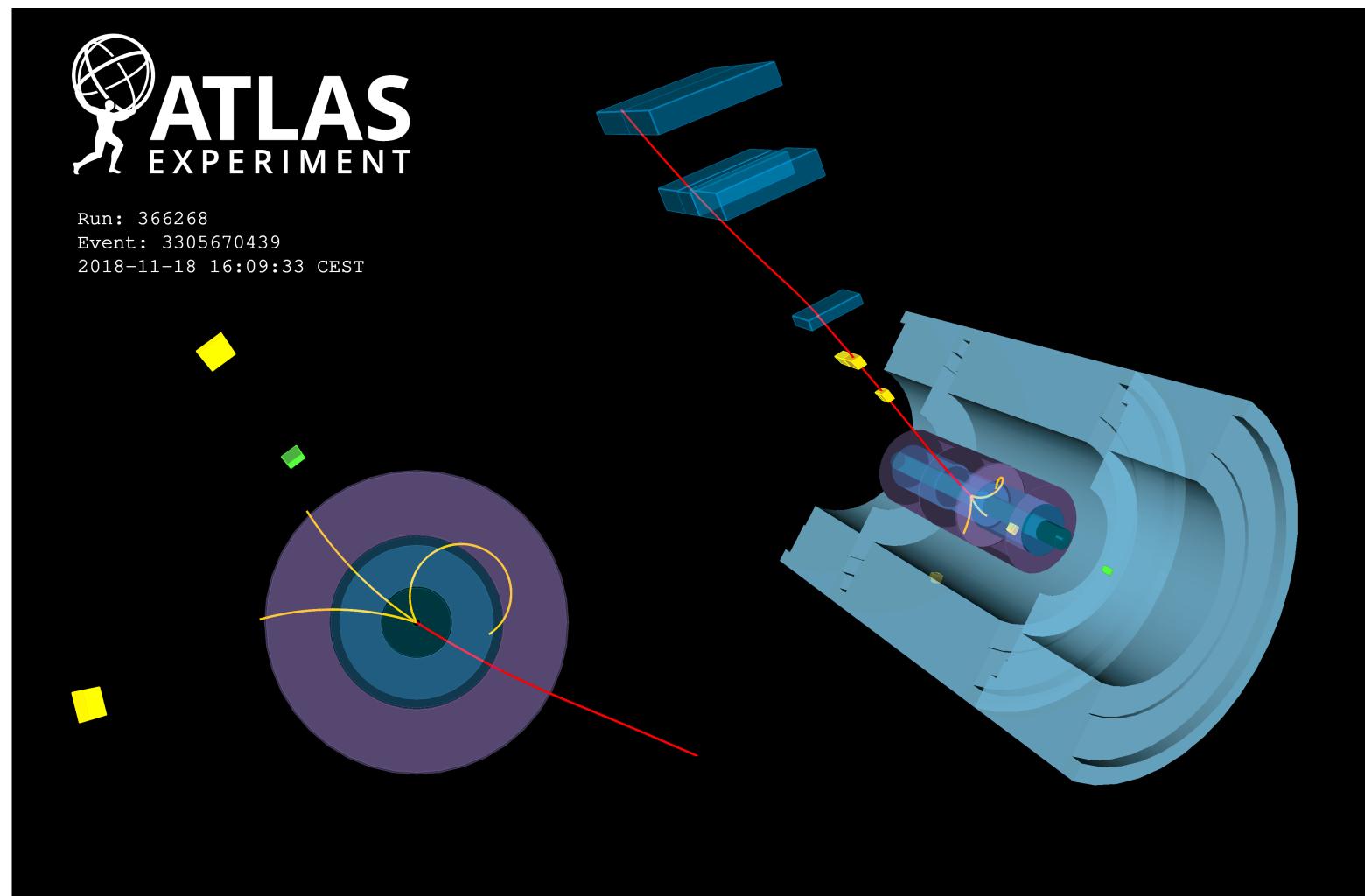
- LHC is the most powerful particle accelerator in the world
 - Four main experiments placed at the LHC ring: ALICE, ATLAS, CMS, LHCb



Recording tau pairs in the ATLAS experiment

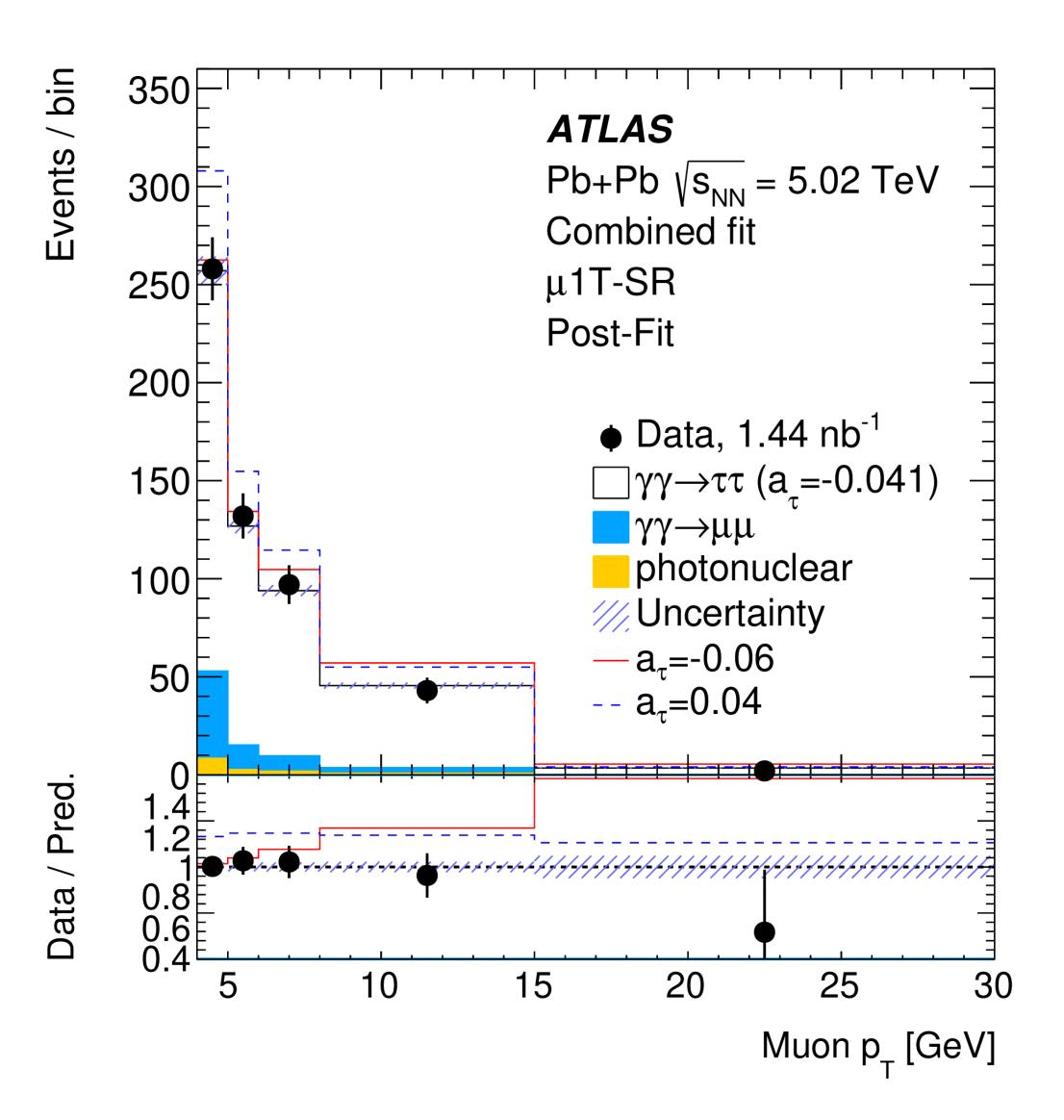
- In ATLAS, we reconstruct decay products of taus
- We observe taus for the first time in ion collisions!





Tau g-2 measured in ATLAS

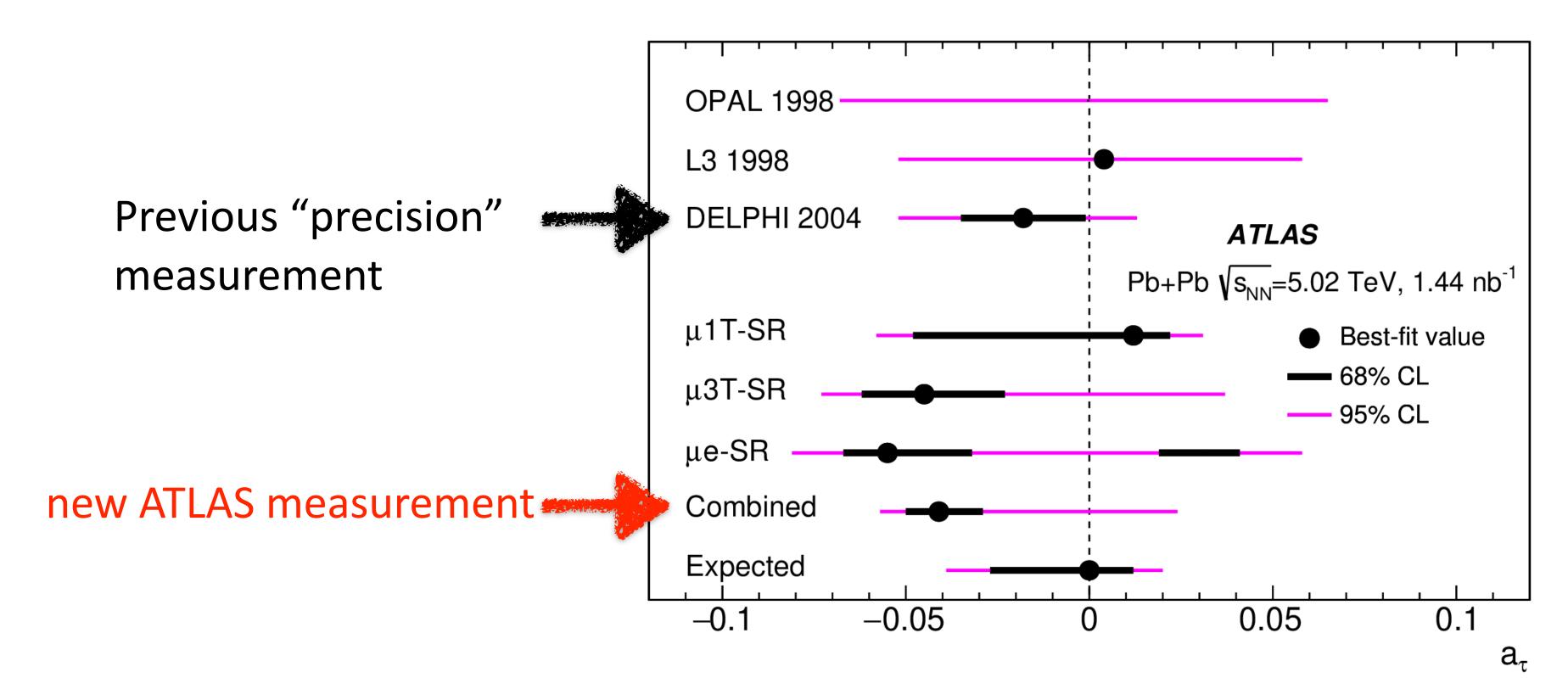
• The value of $a_{\tau}=(g_{\tau}-2)/2$ is sensitive to both cross-section variations and shapes of kinematic distributions



Tau g-2 measured in ATLAS

• We measure $a_{\tau} = (g_{\tau} - 2)/2$

- Phys.Rev.Lett. 131 (2023) 15, 151802
- g_{τ} is found to be consistent with "2" (1.94< g_{τ} < 2.02 @95% CL)
- Note there is also an alternative measurement from CMS experiment (but far less precise)



Summary

- Magnetic dipole moments of charged leptons can be measured with high precision
 - Sensitive to 'new' particles via quantum fluctuations
 - Unprecedented accuracy achieved for electrons and muons: g_e = 2.0023193043612(3), g_μ = 2.0023318411(5)
 - Waiting for a clarification (of the theory): discrepant α measurements (electrons), hadronic contributions (muons)
- Challenging to measure tau lepton magnetic dipole moment
 - ATLAS experiment has measured this recently (for the first time by using heavy-ion collisions) (1.94< g_{τ} < 2.02 @95% CL) -> precision will be improved by studying more data
 - Major involvement of WFiIS staff and students in the preparation of the ATLAS result!

Backup

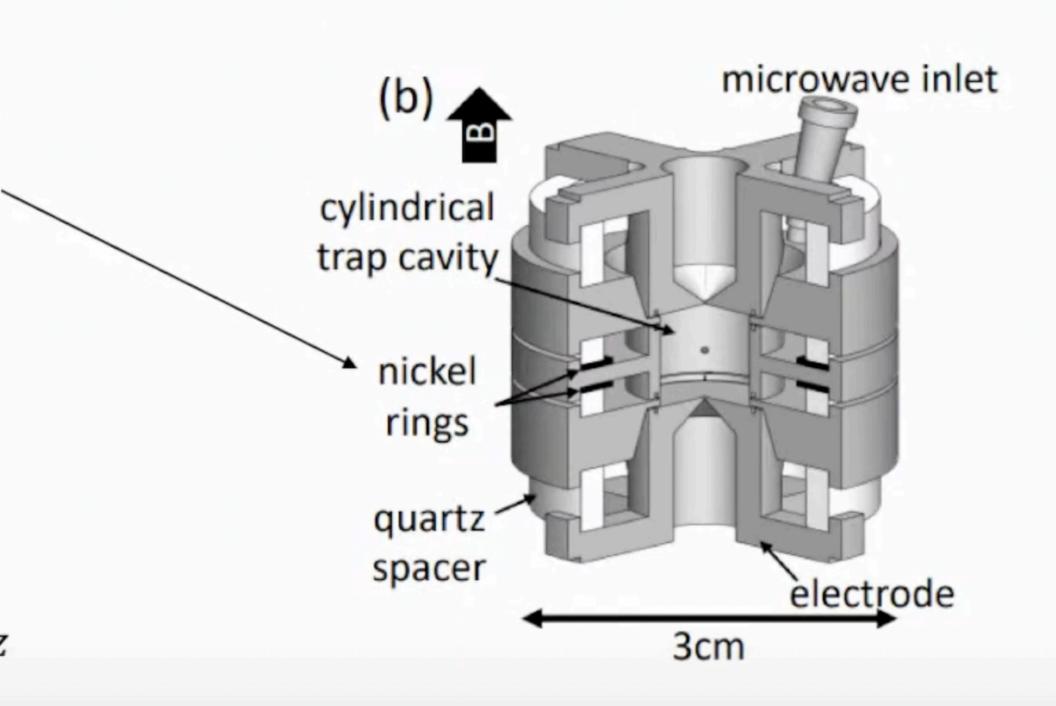
Electron g-2: quantum non-demolition (QND)

Add a small magnetic gradient to couple the spin and cyclotron magnetic moments to the axial frequency $\Delta \vec{B} = B_2 z^2$

Hamiltonian for the axial motion

$$H = \frac{1}{2}m\omega_z^2 z^2 - \mu B_2 z^2$$

$$\downarrow \text{shifts observed } \omega_z$$



QND \rightarrow makes quantum transitions without causing them

$$\Delta \bar{\nu}_z \approx 1.3 \left(n + m_s \right) \text{Hz}$$

Electron g-2: trap imperfections

Imperfect Trap

B in Free Space

Perfect Electrostatic Quadrupole Trap

- tilted B
- harmonic distortions to the trapping potential

$$v_c = \frac{eB}{2\pi m}$$

$$v_c < v_c$$

$$v_z = v_c'$$

$$v_m = v_z$$

$$v_s = \frac{g}{2}v_c$$

$$v_s = \frac{g}{2}v_c$$

$$v_c < v_c$$

$$v_c = \frac{v_c}{v_c}$$

$$v_s = \frac{g}{2}v_c$$

not a measurable eigenfrequency in an Problem: $\frac{g}{}$ imperfect Penning trap

Solution: Brown-Gabrielse Invariance Theorem

$$v_c = \sqrt{(\overline{v}_c)^2 + (\overline{v}_z)^2 + (\overline{v}_m)^2}$$

Electron predictions

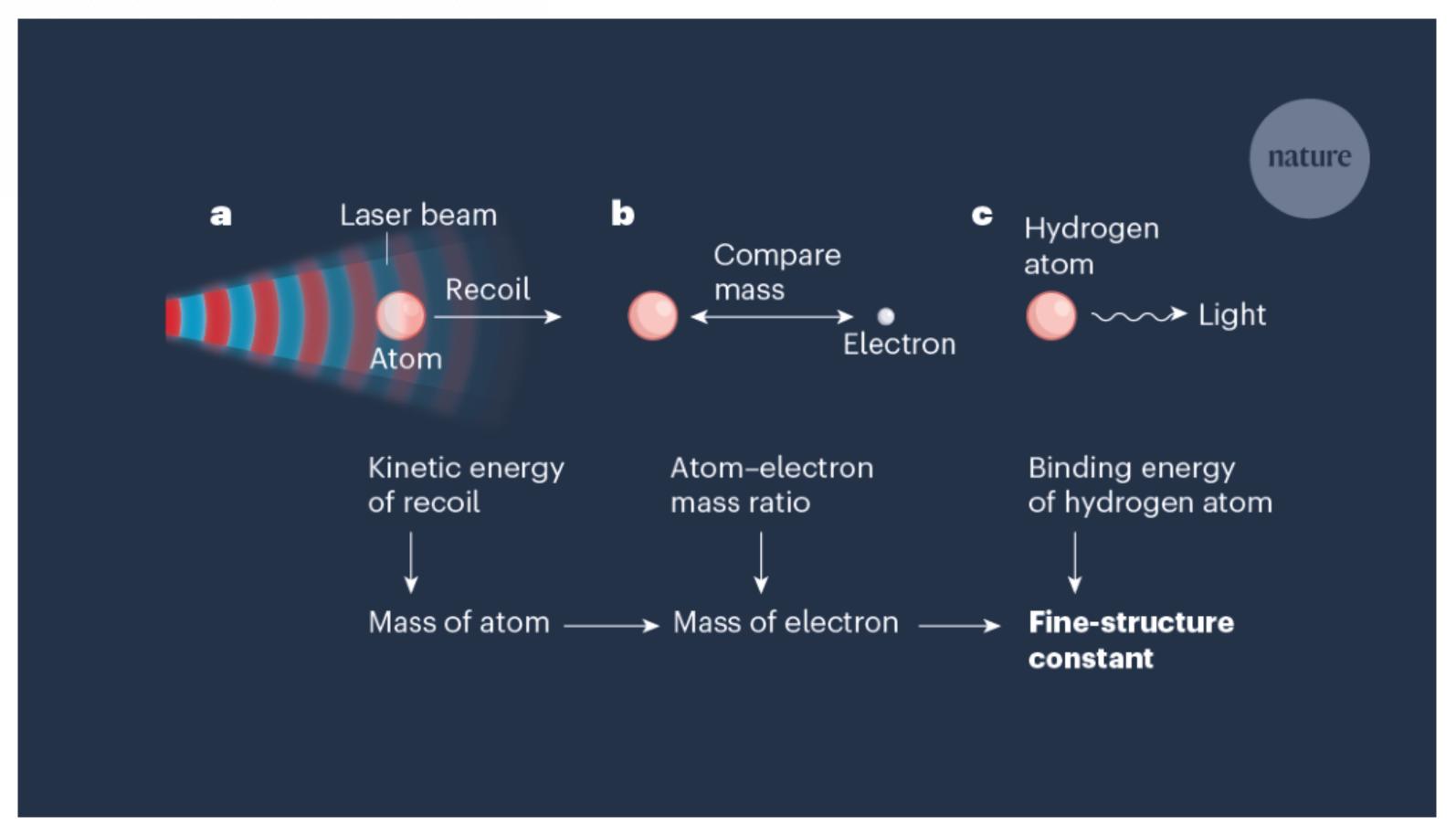
contribution
1.000 000 000 000
0.001 161 409 731 851 (000)(093)
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0.000 000 000 000 456 (011)(000)
0.000 000 000 002 748 (000)
0.000 000 000 001 693 (012)
0.000 000 000 000 031 (000)
1.001 159 652 180 252 (011)(012)(093)

X. Fan, PhD thesis

Fine structure constant measurement

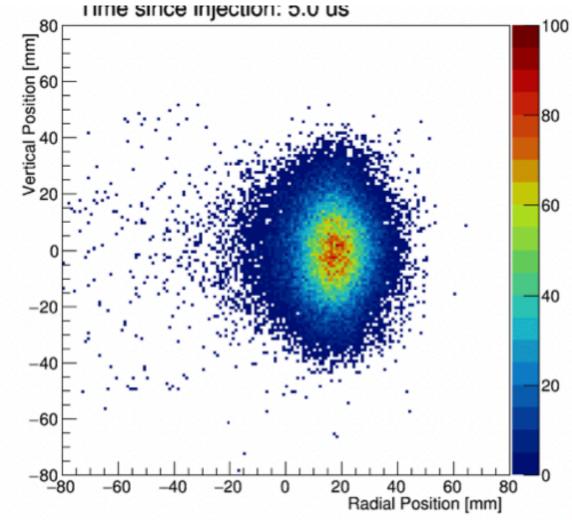
successful independent approach is based on the measurement of the recoil velocity $(v_r = \hbar k/m)$ of an atom of mass m that absorbs a photon of momentum $\hbar k$ (refs. $\frac{10}{11}$). Here \hbar is the reduced Planck constant $(\hbar = h/(2\pi))$ and $k = 2\pi/\lambda$ is the photon wave vector, where λ is the laser wavelength. Such a measurement yields the ratio h/m and then α via the relation

$$lpha^2 = rac{2R_\infty}{c} imes rac{m}{m_{
m e}} imes rac{h}{m}.$$



Muon g-2: beam-dynamics effects

- The muon beam oscillates and breathes as a whole
- The full equation is more complex and corrections due to radial (x) and vertical (y) beam motion are needed



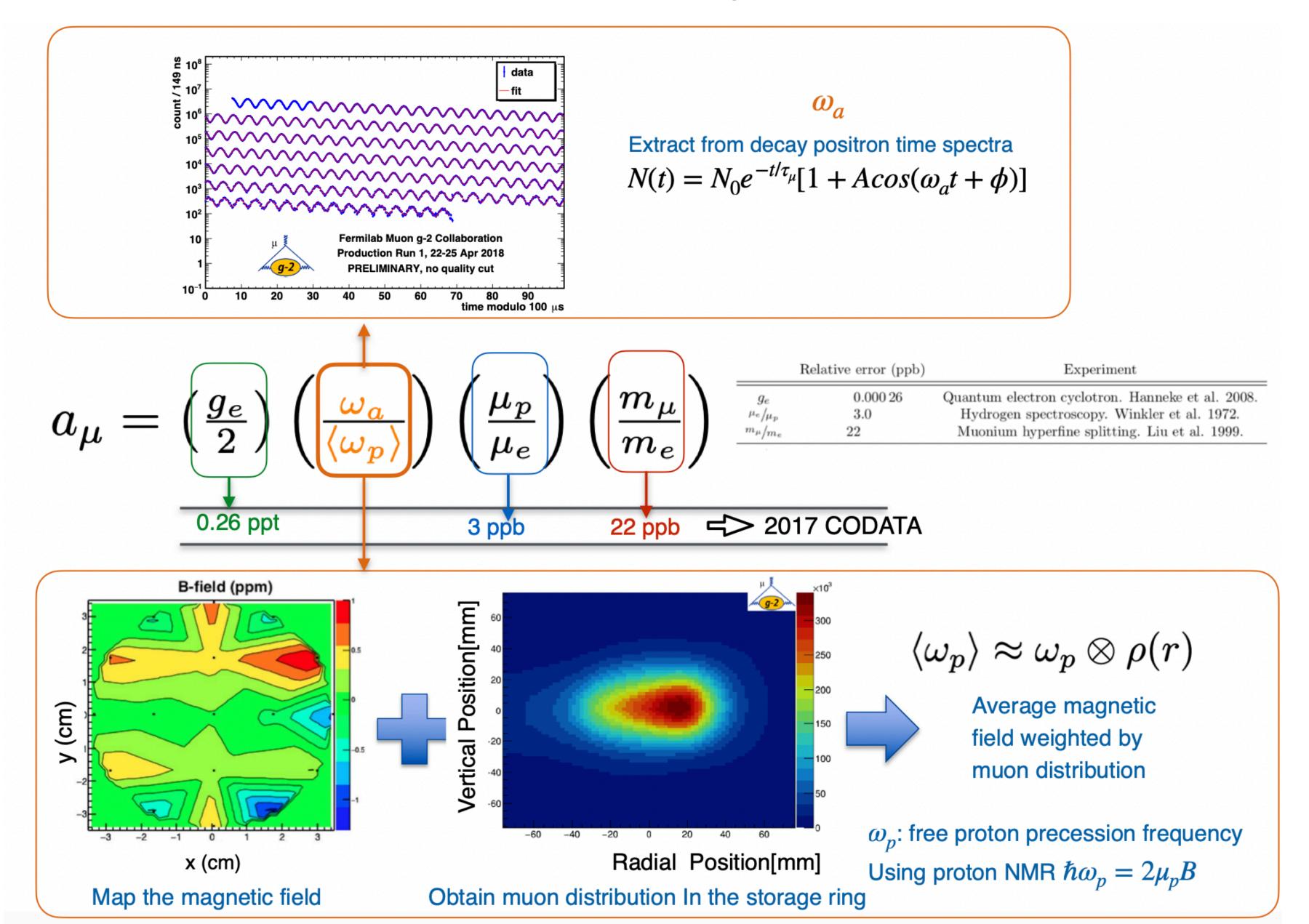
$$\vec{\omega}_{a} = \vec{\omega}_{S} - \vec{\omega}_{C} =$$

$$= -\frac{e}{mc} \left[a_{\mu} \vec{B} - \left(a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} - a_{\mu} \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

- Running at γ_{magic}=29.3 (p=3.094 GeV/c) this coefficient is null
- Because of momentum spread (<0.2%) ->
 E-field Correction

Vertical beam
 oscillation →
 Pitch correction

Muon g-2: extracting a_µ



Tau g-2: a_tau parametrisation

 Elementary γγ→τ τ cross section has explicit dependence on photon-τ vertex function:

$$i\Gamma_{\mu}^{(\gamma\ell\ell)}(p',p) = -ie\left[\gamma_{\mu}F_{1}(q^{2}) + \frac{i}{2m_{\ell}}\sigma_{\mu\nu}\underline{q^{\nu}}F_{2}(q^{2}) + \frac{1}{2m_{\ell}}\gamma^{5}\sigma_{\mu\nu}\underline{q^{\nu}}F_{3}(q^{2})\right]$$

$$= a_{\tau} (q^{2}=0) \qquad = d_{\tau}*2m_{\tau}/e (q^{2}=0)$$