

Akademia Górniczo-Hutnicza im. Stanisława Staszica w Krakowie

Laser rentgenowski na swobodnych elektronachnowe źródła promieniowania X w DESY-HASYLAB

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Light sources for research with photons at DESY



Total project construction cost - 1,081.6 M€





Petra III Facilities

a total investment of 225 million $\ensuremath{ \in }$

PETRA and HERA accelerators



The development of X-ray radiation sources



Average and Peak Brilliance in Comparison



The advantages of XFEL

•ultra-high brilliance and photon flux intensity

- •high collimation
- •wide energy/wavelength tunability (up to hard X-rays)
- specific polarization
- •spatial coherence (transverse and longitudinal)
- pulsed time structure

Free-electron laser and conventional laser

LASER-Light Amplification by Stimulated Emission of Radiation

XFEL- X-Ray Free Electron Laser based on SASE- Self-Amplified Spontaneous Emission



Principle of a conventional quantum laser where the electrons are bound to atomic, molecular or solid-state energy levels ("bound-electron laser").

Plasma X-ray laser



Figure 2. Rendering of Livermore's COMET (compact multipulse terawatt) tabletop x-ray laser showing the laser system and target chamber. The inset shows laser beams hitting the stepped target and producing a plasma, which in turn generates an x-ray laser beam.

The FEL is not, strictly speaking, a laser, i.e. a device based on quantummechanical stimulated emission, and its operation is completely described within the framework of classical physics. The FEL is a system consisting of a relativistic electron beam and a radiation field interacting with each other while propagating through an undulator.





Schematic representation of undulator radiation.

XFEL_Linac_Accelerator.mov

XFEL_SASEFEL.mov



Free-electron laser configurations: oscillator (top), seeded amplifier (middle), and SASE FEL (bottom).



Principle of a free-electron laser: (a) For visible or infrared light (b) In the ultraviolet and X-ray region one can apply the mechanism of self-amplified spontaneous emission



The self-seeding mechanism:

Wavelength of undulator radiation

$$v_z = c \left(1 - \frac{2 + K^2}{4\gamma^2} \right)$$

Relativistic Doppler shift

$$\lambda_t = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} \right)$$

$$K = \frac{eB_0\lambda_u}{2\pi m_0 c} \approx 1$$

B₀- peak magnetic field in the undulator

Source: DESY, Hamburg

v<c - electron velocity

v_z - electron velocity component along the z axis γ - Lorentz factor

$$\gamma = \frac{m}{m_0} = \frac{W}{W_0}$$

for $\lambda_u = 30 \text{ mm i } \gamma = 1000$

$$\lambda^{*} = \frac{\lambda_{u}}{\gamma} = 30 \mu m$$
$$\lambda_{t} \approx \frac{\lambda^{*}}{2\gamma} = \frac{\lambda_{u}}{2\gamma^{2}} = 15 nm$$
source:



Condition for sustained energy transfer from electron to light wave in an undulator: The light wave has to slip forward by per half-period of the electron trajectory. Plane electromagnetic wave

$$E_{x}(z,t) = E_{0} \cos(k_{t}z - \omega_{t}t)$$

with $k_{t} = \omega_{t} / c = 2\pi / \lambda_{t}$

The time derivative of the electron energy

$$W = \gamma m_0 c^2$$
 is

$$\frac{dW}{dt} = \vec{v} \bullet \vec{F} = -ev_x(t)E_x(t)$$

$$\lambda_{t} = \frac{\lambda_{u}}{2\gamma^{2}} \left(1 + \frac{K^{2}}{2} \right)$$

To achieve light amplification the electron energy must exceed the resonance energy

$$W > W_t \equiv \gamma_r m_0 c^2$$

The resonant Lorentz factor is

$$\gamma_r = \sqrt{\frac{\lambda_u}{2\lambda_t}} \left(1 + \frac{K^2}{2} \right)$$





microbunching.mov

The exponential growth of the FEL pulse energy E as a function of the length z traveled in the undulator.



Average energy in the radiation pulse against magnetic undulator length for SASE 3 operating at 0.4 nm

microbunching.mov

Numerical simulation of microbunching

The power gain

$$P(z) = \frac{P_{in}}{9} \exp(2z/L_g) \text{ for } z \ge 2L_g$$

L_g- field gain lenght P_{in}- effective input power

$$L_{sat} = 10 \bullet L_g$$
 $L_g = \frac{\lambda_u}{4\pi\rho}$

 ρ -saturation efficiency L_{sat} - saturation lenght

Electron energy – 17.5 GeV Bunch peak current- 5 kA Peak power- 100 TW Conversion efficiency-0.1% Peak power of X-rays- multiGW X-ray intensities- 10¹⁸ W/cm²

Location of the European XFEL

Cut through the electron source of FLASH.

- **Requirements on the drive beam**
- of an X-ray FEL
- •High peak current
- •Very low emittance
- •Small energy spread

Principle of longitudinal electron bunch compression.

Electron bunch time pattern with 10 Hz repetition rate and up to 3,000 bunches in a 0.6 ms long bunch train. The separation of electron bunches within a train is 200 ns for full loading. The duration of electron bunches is ~200 fs and the nonlinear FEL process reduces the duration of the photon pulses to ~100 fs.

Schematic layout of the planned X-ray laser laboratory in a top view. Black and red lines indicate electron and photon beam lines, respectively. The undulators are marked in blue and violet.

	Photon energy [keV]	Polarisation	Tunability	Gap varia-
tion				
SASE 1	12.4	Linear	No	Yes
SASE 2	3.1 – 12.4	Linear	Yes	Yes
SASE 3	0.8-3.1 (0.25-1.0)*	Circular/Linear	Yes	Yes
U1, U2	20 - 100	Linear	Yes	Yes

Scientific applications of XFEL radiation

- •Femtosecond chemistry
- Investigation of single molecules
- X-ray microscopy
- Plasma physics
- Coherent scattering experiments

Femtochemistry

"Filming" chemical reactions using ultra-fast lasers. The brilliance of one single laser flash is so high that it can generate images of reacting molecules with atomic resolution Source: DESY, Hamburg

X-ray microscopy

Ribosomes

Source: DESY, Hamburg

Structural biology

> The X-ray laser opens up completely new opportunities to decipher biological molecules with atomic resolution without the need for the extra step of growing crystals. The X-ray laser flashes are so intense that they can be used to create a high-resolution image of a single molecular complex. The flash duration is shorter than 100 femtoseconds and is thus short enough to produce an image before the sample is destroyed by the intense X-rays.

> Ribosomes are large molecular complexes that act as "protein factories" and occur in every cell. The X-ray laser opens up completely new opportunities to decipher such biological structures with atomic resolution without the need for the extra step of tediously growing them into crystals first.

Diffraction imaging

Schematic diagram of the single-particle diffraction imaging experiment.

Diffraction of proteins

Planar section through the centre of the molecular transform of a small protein molecule (lysozyme, left) and of a larger virus capsid (tomato bushy stunt virus, right)

Biomolecules are destroyed by intense X-ray radiation: they "explode". The illustration shows a simulation of this process. In order to obtain a usable image of the biomolecule, the image must be recorded very quickly before the sample is destroyed.

Femtochemistry

A chemical reaction is triggered by a laser flash: The molecules are excited to a more energetic state, after which they "drop" back to their less energetic equilibrium state. A second laser pulse is then sent at varying intervals (t) after the first one to take instantaneous snapshots of the changes that have occurred in the molecule.

Plasma physics

Experimental arrangement for spectroscopy experiments. The burial of a microdot in a matrix helps to better define the plasma conditions at the time and plac of excitation by the XFEL pump. Time- and space-resolved spectroscopy will be used todiagnose the emission.

Figure 6.4.12 Simulation of change to emission spectrum for an Aluminium plasma pumped with an XFEL tuned to 1,869 eV in order to pump the Helium-like 1s²-1s3p transition.

PETRA III

PETRA III

Number	Name		Energy range	Comment	Contact
P01	Dynamics beamline	Dynamik Strahlführung	5 - 40 keV		H. Franz, DESY
P02	Hard X-ray scattering	Harte Röntgenstrahlen/Beugung	20 - 100 keV	straight	DESY
P03	Micro SAXS/WAXS	Mikro SAXS/WAXS	8 - 25 keV	down	S. Roth, DESY
P04	Variable Polarization XUV	XUV Strahlführung	0.2 - 3.0 ke∨		J. Viefhaus, DESY
P05	Micro- and nano-tomography	Mikro/Nano-Tomographie	8 - 25 keV	side	A. Haibel, DESY/GKSS
P06	Hard X-ray micro probe, imaging	Nano-Analyse/Abbildung	2.4 - 50 keV	straight	G. Falkenberg, DESY
P07	High energy materials science	Materialforschung	40 - 300 ke∨		N. Schell, GKSS/DESY
P08	High resolution diffraction	Hochauflösende Diffraktion	5.4 - 30 keV	top	O. Seeck, DESY
P09	Resonant scattering/diffraction	Resonante Streuung/Diffraktion	2.4 - 50 keV	straight	J. Strempfer, DESY
P10	Coherence applications	Kohärenzanwendung	4 - 25 keV		O. Leupold, DESY
P11	Bio imaging/diffraction	Proteinkristallographie/Abbildung	8 - 25 keV	side	MPI, HGF, DESY
P12	BioSAXS	Bio SAXS	4 - 20 keV	straight	M. Rößle, EMBL
P13	Macro molecular crystallography I	Biokristallographie I	5 - 35 keV	side	M. Cianci, EMBL
P14	Macro molecular crystallography II	Biokristallographie II	5 - 35 keV	straight	G. Bourenkov, EMBL

high beta section	142x5 μm
low beta section	35x6 µm

Positron energy 6 GeV 14 beamlines with 30 experimental stations X-ray spectral range from 0.2 to 300 keV, t=1ps Source: DESY, Hamburg