

# **Renormalized Wannier functions at the border of Mott localization**

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## **Collaboration:**

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**Robert Podsiadły – Jag. Univ., Krakow**

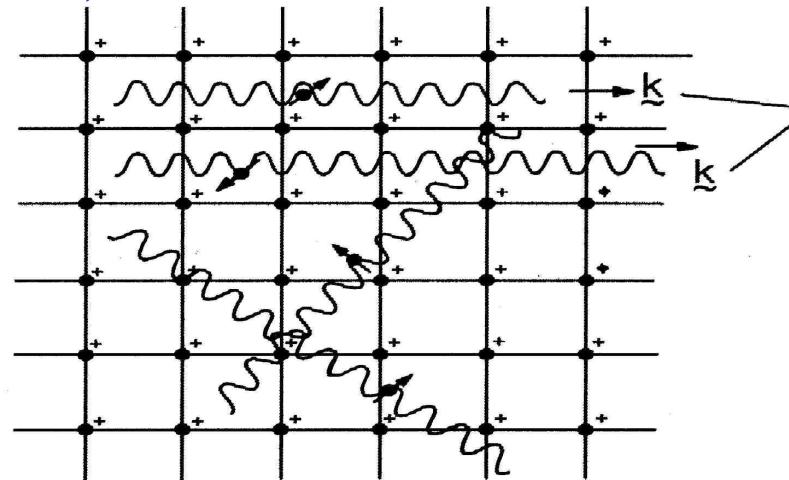
**Włodek Wójcik – Tech. Univ., Krakow**

# **Plan**

- 1. From atoms to metals**
- 2. Wave function readjustment in the correlated electron state**
- 3. Example: exact solution for nanosystems and Hubbard chain**
- 4. Quantum critical behavior of the wave function (size)**

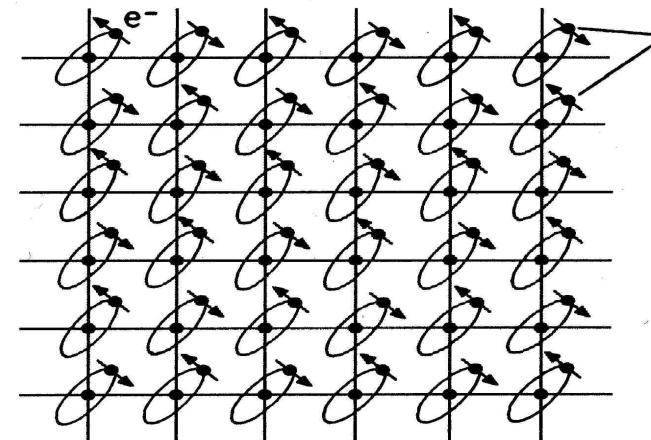
# Delocalized versus localized

a) Metal

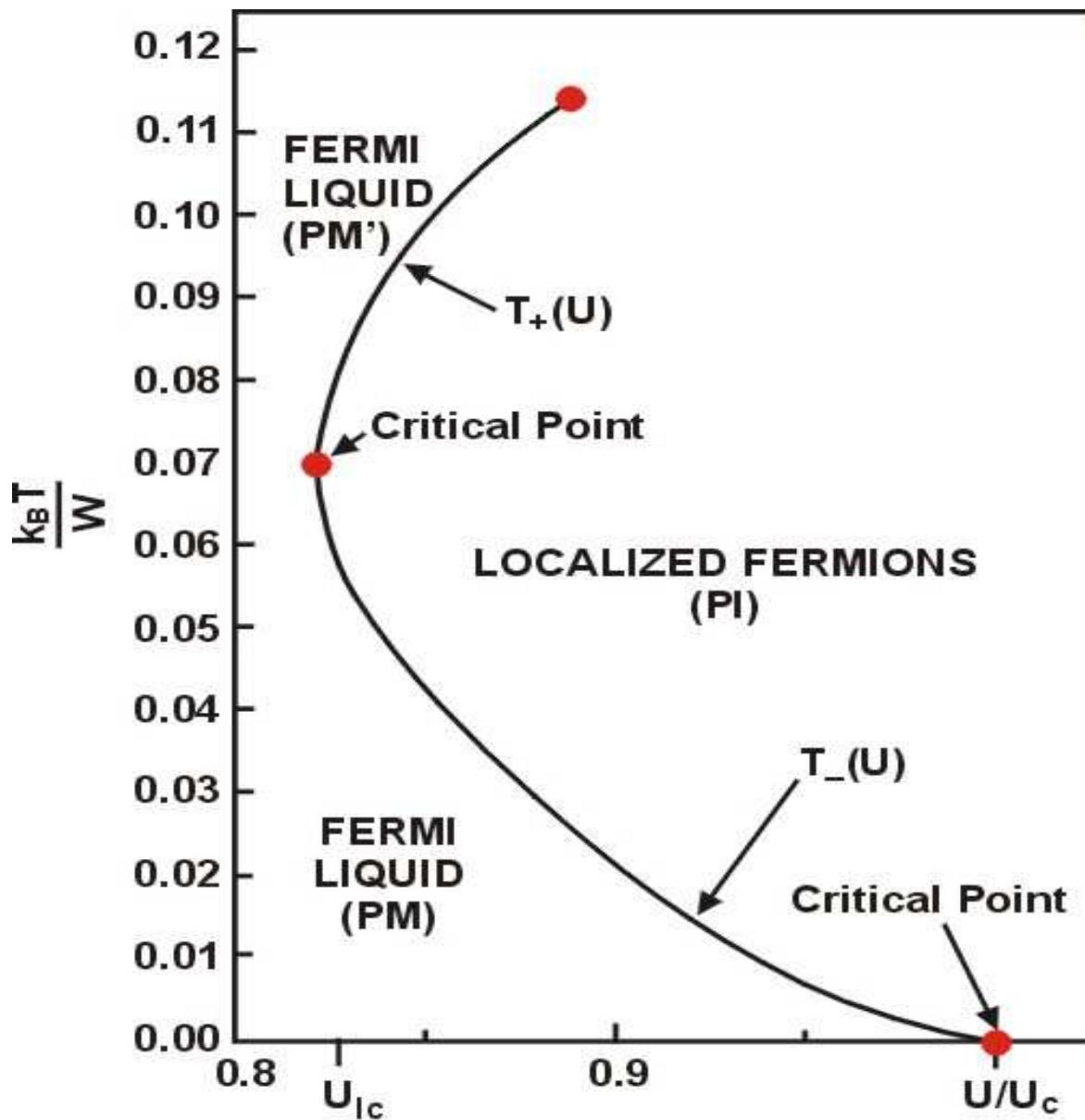


Plane waves  
(Bloch states)

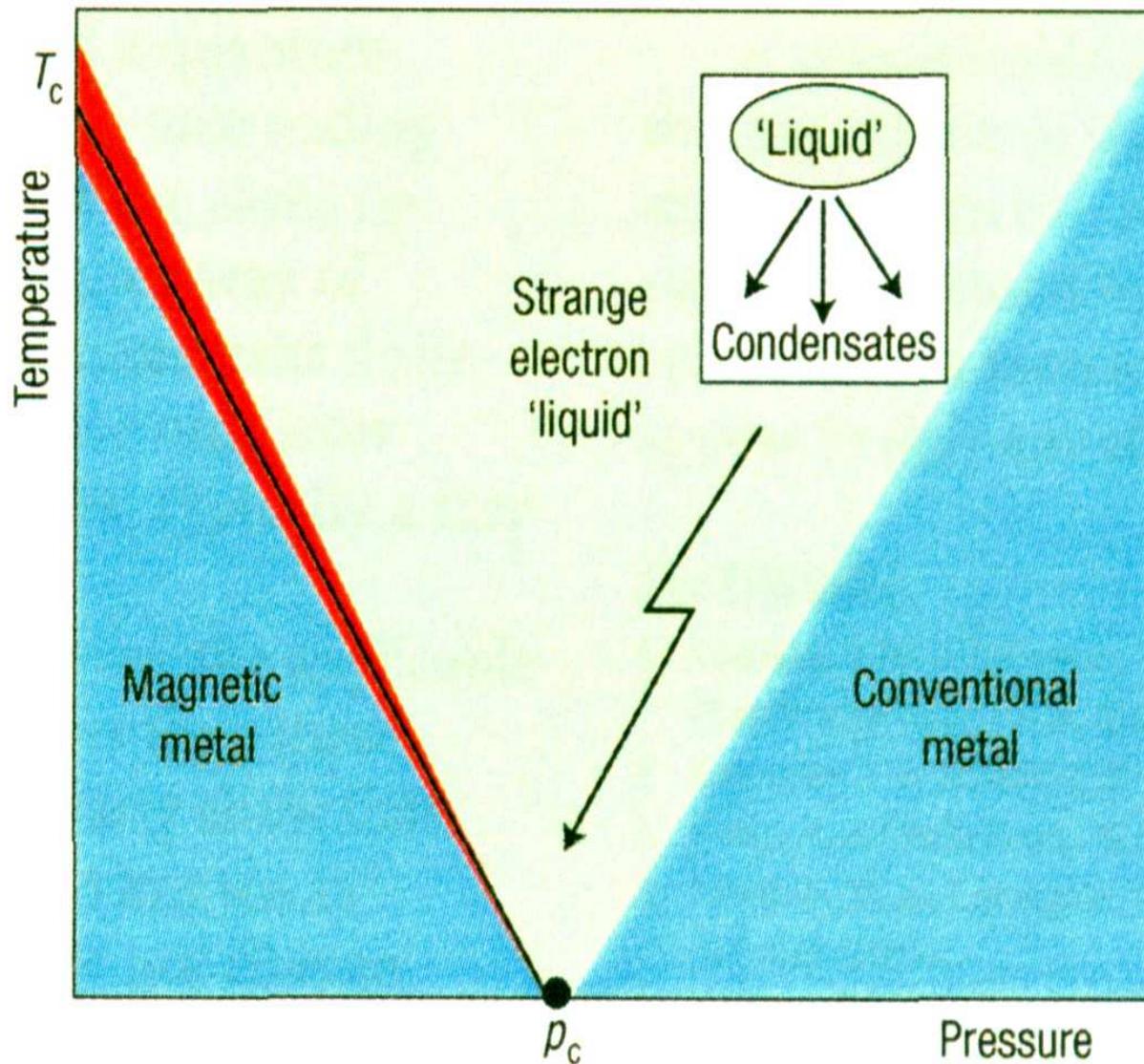
b) Mott-Hubbard insulator



Atomic states

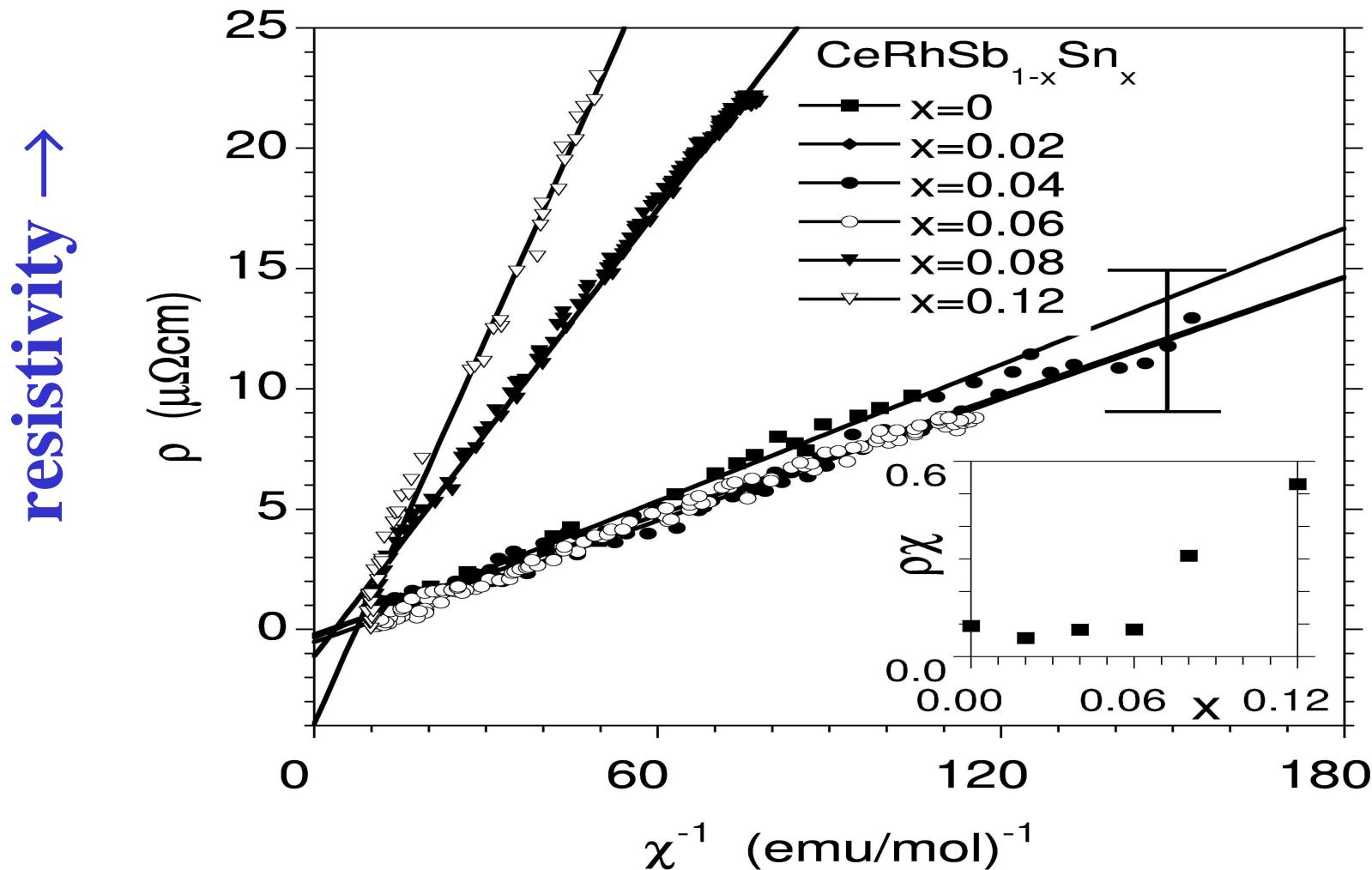


J. S. et al., PRL **59**, 728 (1987) – orbitally nondegenerate;  
A. Klejnberg & J. S., PRB **57**, 12 041 (1998) – degenerate.

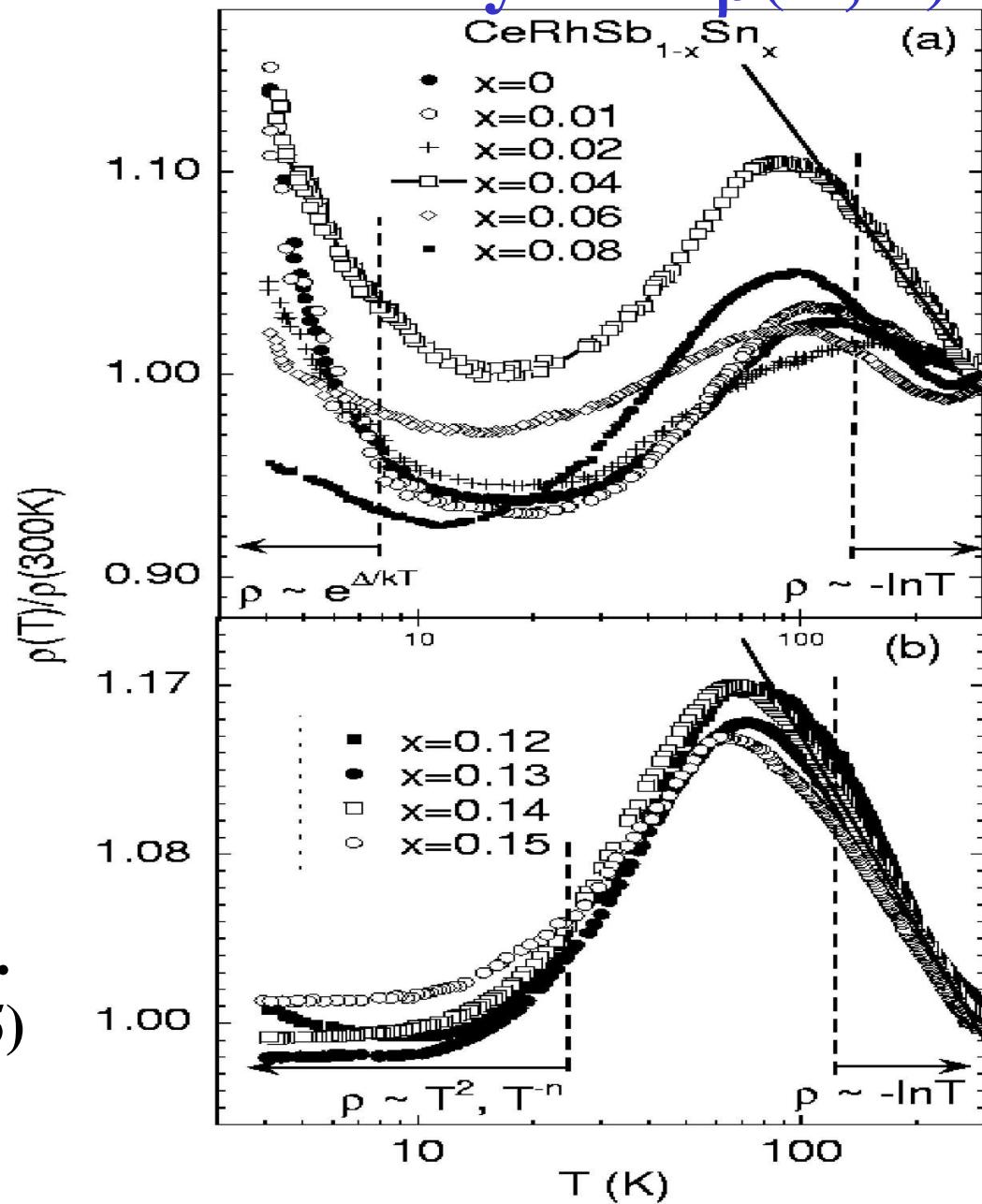


G. Lonzarich, Nature (2005)

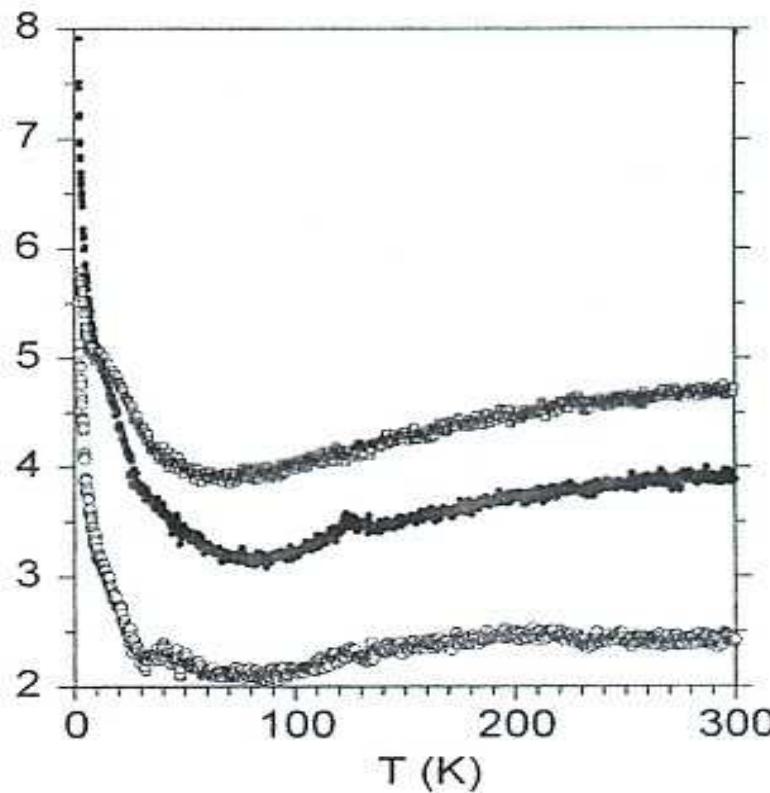
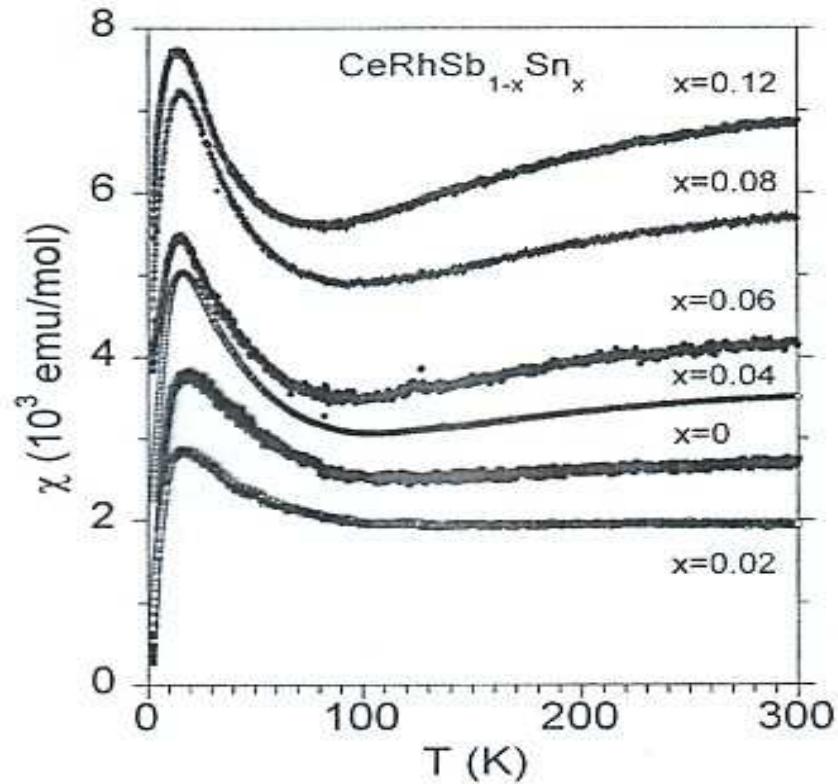
# Universal scaling



# Resistivity data $\rho(T; x)$



A.Š. & J.S.  
PRL (2005)



cond. el.

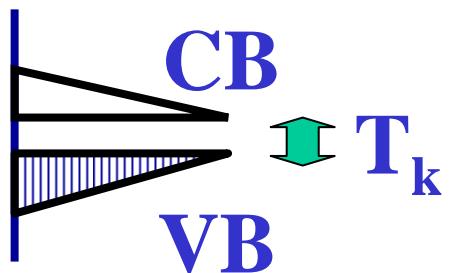
$S_{ce}$

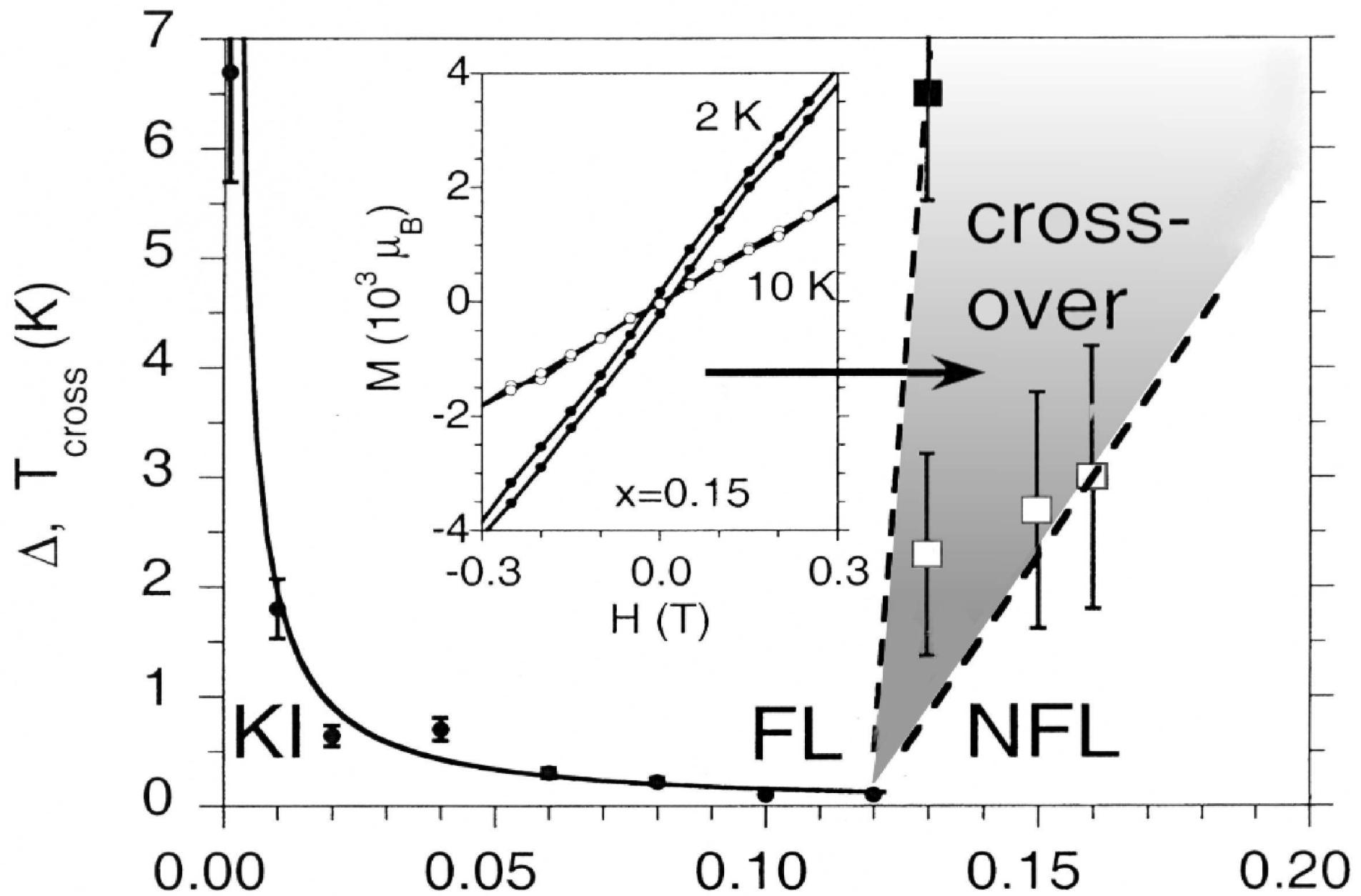
$S_{tot} \equiv 0$

or  
else:

↓ ↓ ↓ ↓

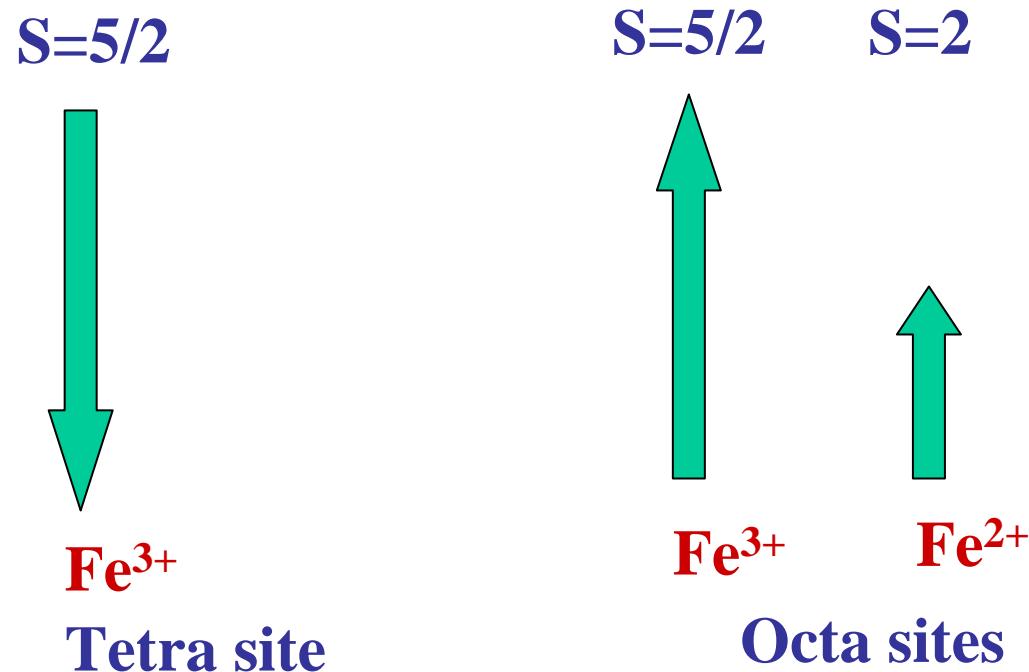
↑



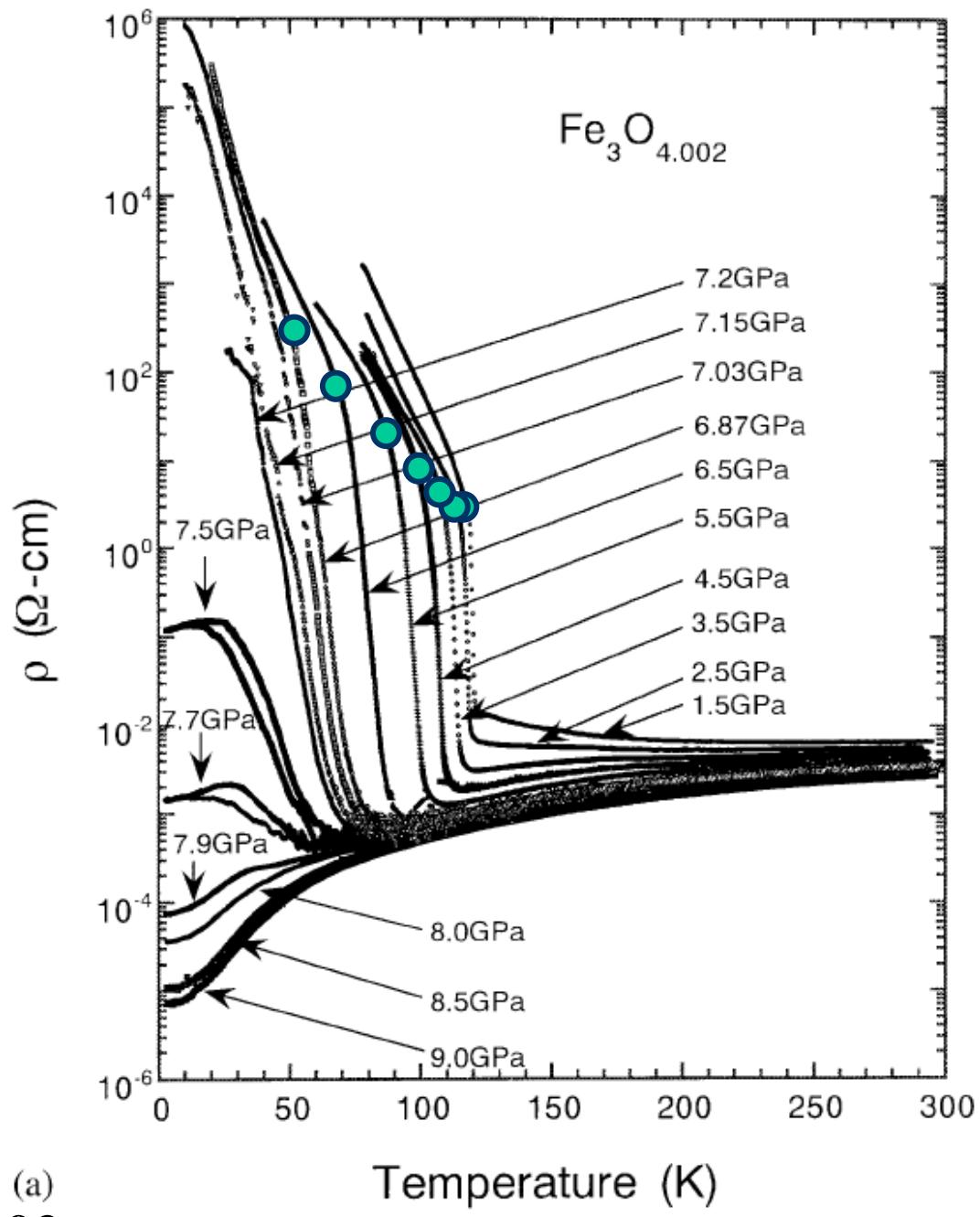


## II. Metallization of magnetite

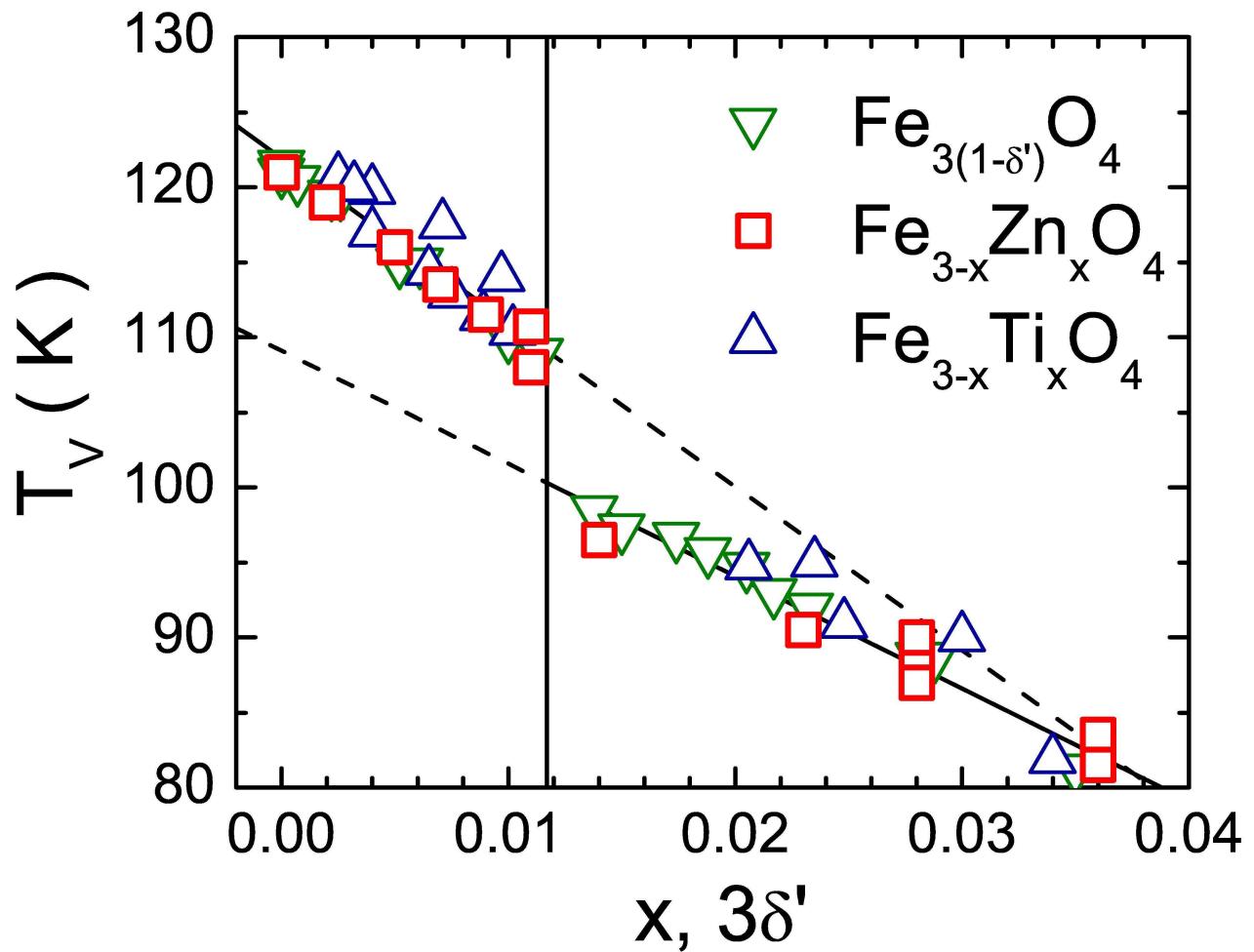
- Ferrimagnetic material:  
 $T_c=860 \text{ K}$ ,  $M=4.1 \text{ Bohr magnetons}$



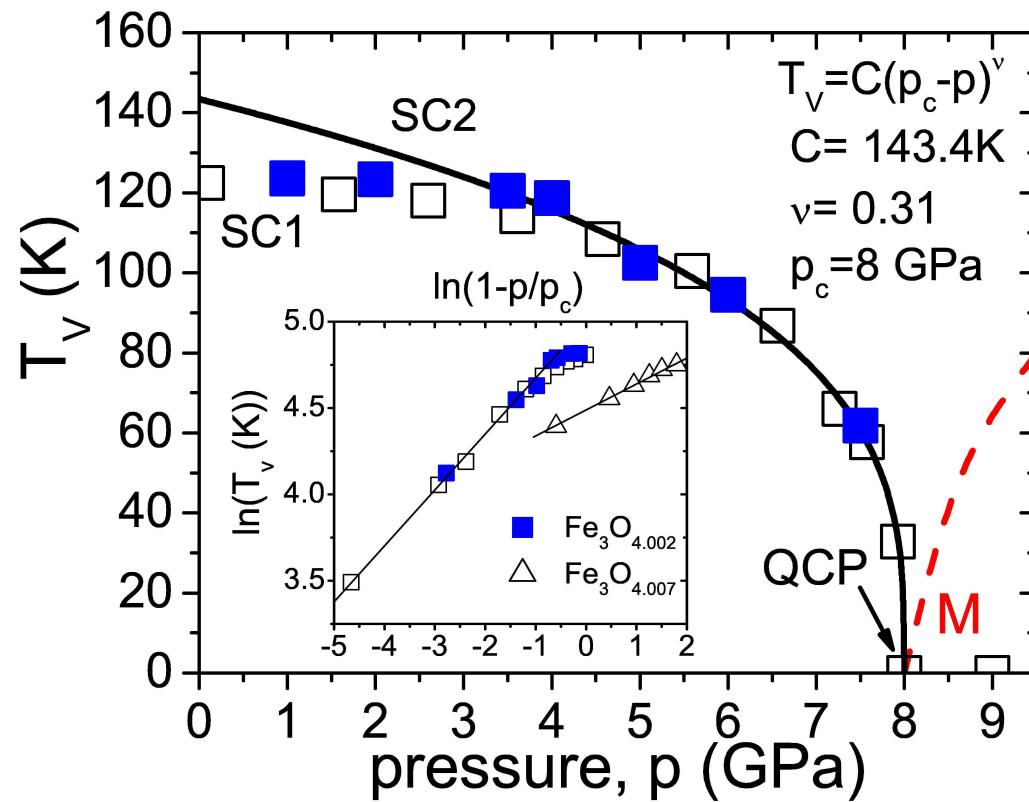
- Verwey transition:  $T_V = 122 \text{ K} \pm 1 \text{ K}$   
(at  $p = 0$ )



N. Mori et al. (a)  
Physica B, 2002



Z. Kąkol, A. Kozłowski, Z. Tarnawski,...J.M. Honig



J.S., A. Kozłowski, Z. Tarnawski, Z.Kąkol,,  
 Y.Fukami, F.Ono, R. Zach, L.J. Spałek, and  
 J.M. Honig,  
 Phys. Rev.B 78, 100401 (R) (2008)

# Localization criterion: Mott

Kinetic energy in e<sup>-</sup> gas/particle

$$\overline{\epsilon} = \frac{3}{5} \epsilon_F = \frac{3}{5} \frac{\hbar^2}{2m^*} \left( 3\pi^2 \frac{N}{V} \right)^{2/3} \sim \rho^{2/3}$$

$$\epsilon_{e-e} = \frac{1}{2} \frac{e^2}{\epsilon d_{e-e}} = \frac{e^2}{2\epsilon} \rho^{1/3}$$

$$d_{e-e} = \left( \frac{V}{N} \right)^{1/3}$$

$\overline{\epsilon} = \epsilon_{e-e}$   **gas instability**

$$\underbrace{\left( \frac{\hbar^2}{m^* e^2} \epsilon \right)}_{a_B} \rho_c^{1/3} = \frac{5}{3} \frac{1}{(3\pi^2)^{2/3}} \approx 0.17$$

$$a_B \cdot \rho_c^{1/3} \approx 0.17 \sim 0.2$$

⇒ Fermi - sphere collapse

In one dimension:

$$a_B \rho_C \approx 1 \Rightarrow R_C \approx a_B$$

# **Microscopic many-particle Hamiltonian for a nanosystem**

(for extended system no phase factor in t)

$$H = \epsilon_a^{\text{eff}} \sum_j n_j + t \sum_{j\sigma} \left( e^{-i\phi/N} c_{j\sigma}^\dagger c_{j+1\sigma} + \text{h.c.} \right) + U \sum_i n_{i\uparrow} n_{i\downarrow} + \sum_{i < j} K_{ij} \delta n_i \delta n_j,$$

$$\delta n_i \equiv n_i - 1, \quad \epsilon_a^{\text{eff}} = \epsilon_a + N^{-1} \sum_{i < j} (2/R_{ij} + K_{ij})$$

The microscopic parameters t, U, K should be calculated together with the H diagonalization -> wave function optimization in the correlated state

**The main ingredient:  
interelectronic correlations  
and wave function treated on  
the same footing**

**The main result:  
evolution of the many-atom  
system as a function of  
interatomic spacing**

Single-particle  
Schrödinger eq.

$$\sum_j H_{ij} w_j(\mathbf{r}) = \epsilon_i w_i(\mathbf{r})$$

**EDABI**

Single-particle basis

$$\{w_i(\mathbf{r})\}$$

Field operators

$$\hat{\Psi}(\mathbf{r}), \hat{\Psi}^\dagger(\mathbf{r})$$

Diagonalization  
in the Fock space

$$H = |\Psi_0\rangle E_G \langle \Psi_0| + \dots$$

Ground-state energy

$$E_G = \langle \Psi_0 | H | \Psi_0 \rangle$$

Single-particle  
basis optimization

$$\{w_i^{\text{ren}}(\mathbf{r})\}$$

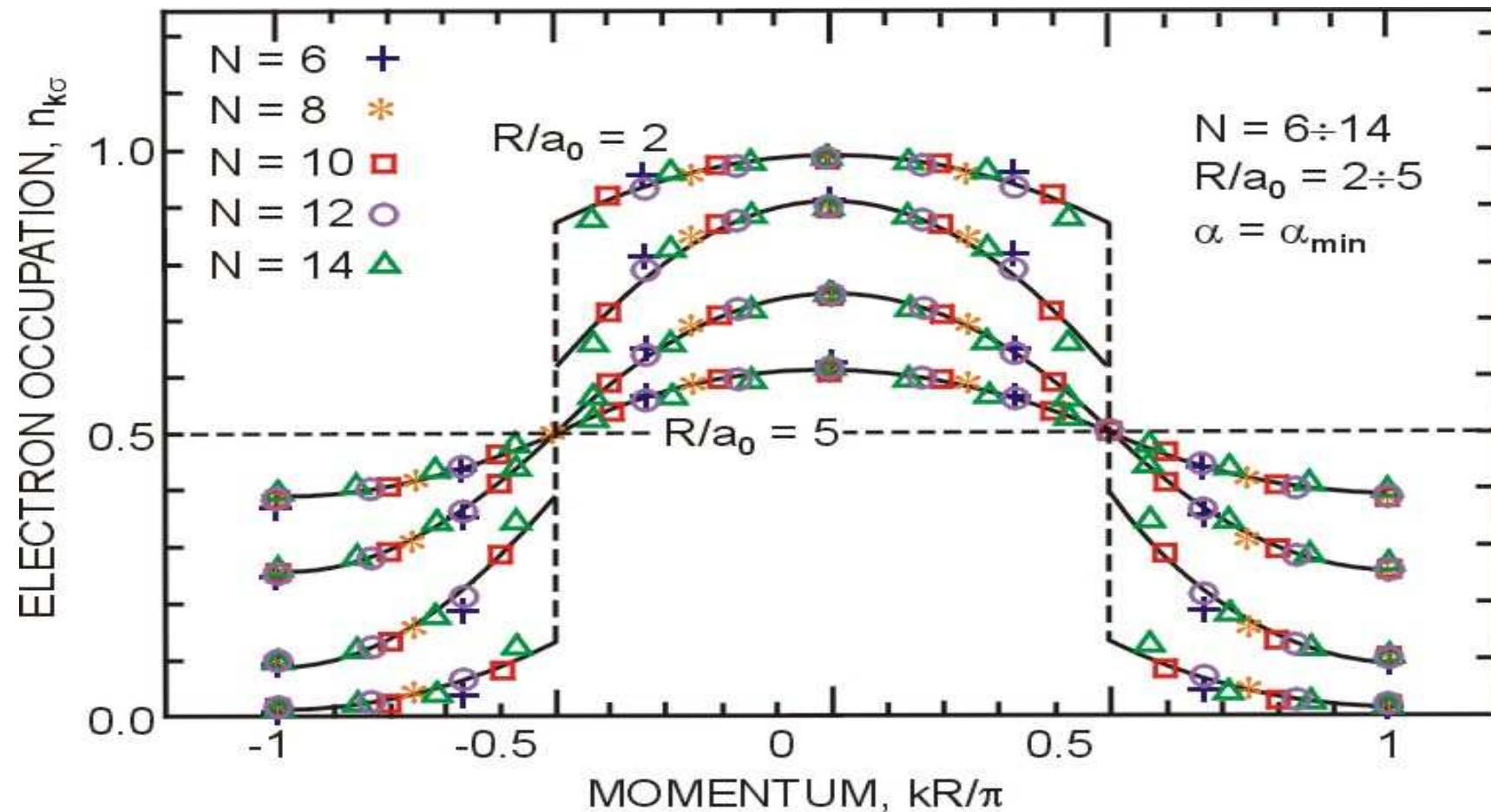
$$\hat{\Psi}^{\text{ren}}(\mathbf{r}), (\hat{\Psi}^{\text{ren}})^\dagger(\mathbf{r})$$

$$\Psi_0^{\text{ren}}(\mathbf{r}_1, \dots, \mathbf{r}_N)$$

*Renormalized N-particle  
wavefunction*

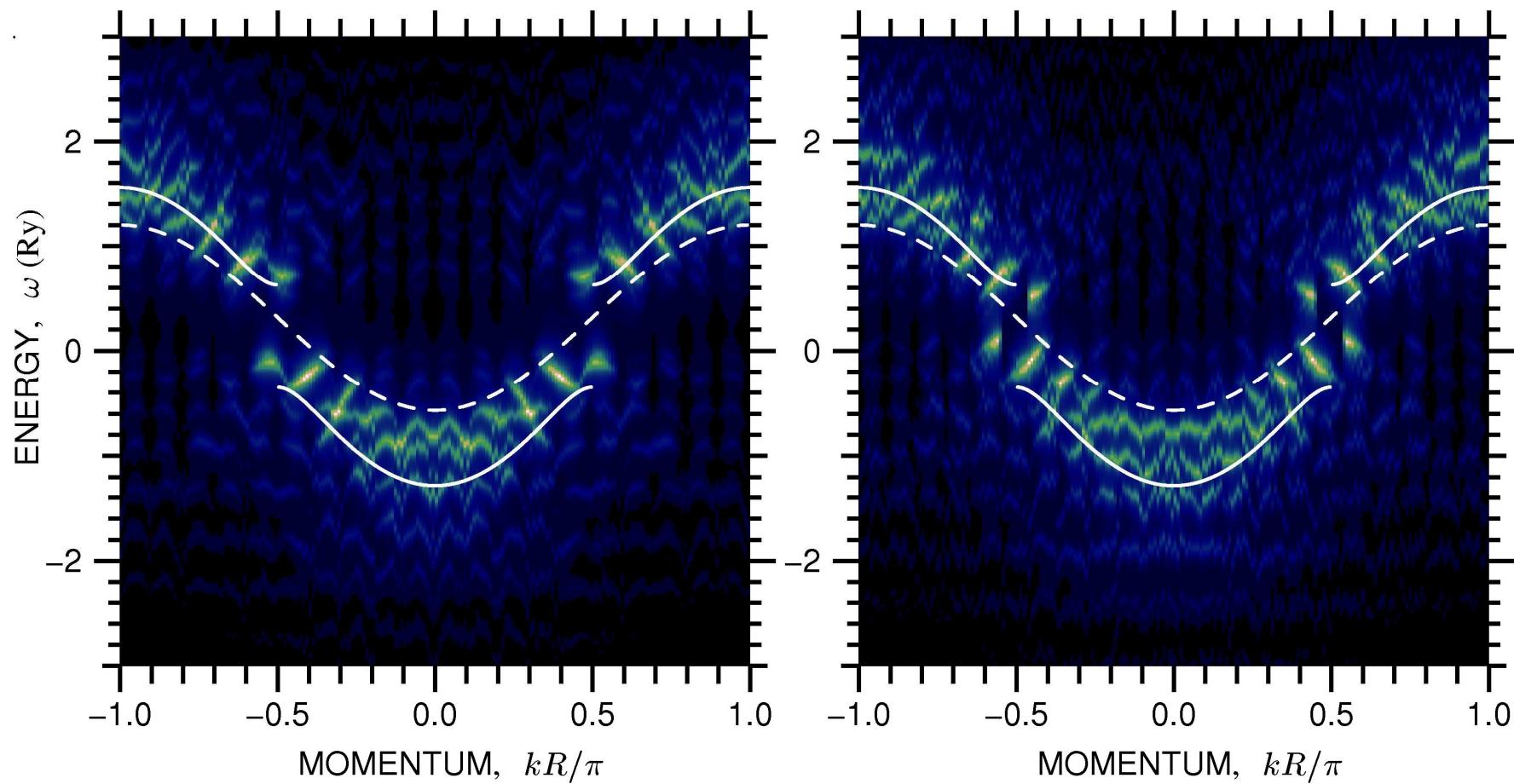
**J. Spałek, R. Podsiadły, W. Wójcik, and A. Rycerz,  
Phys.Rev. B 61, 15676 (2000);PRB (2001-2002)**

# Momentum distribution:Fermi-Dirac vs continuous



J. S. & A. Rycerz, PRB-R (2001-2004); review: 2007;  
Didactical: J.S., in Encyclopedia of Condensed Matter  
Physics, Elsevier, vol. 3, pp. 126-136 (2005)

# **Renormalized band energies: even and odd**



# Extended systems : EDABI

## D=1

# Supplement: Infinite Hubbard chain vs. nanochain

Ground state energy functional:

$$\frac{E}{N} = \epsilon_a^{\text{eff}} - 4t \int_0^\infty \frac{J_0(\omega)J_1(\omega)}{\omega [1 + \exp(\omega U / 2t)]} d\omega$$

$$\delta n_i \equiv 1 - n_i \equiv 0$$



Periodic bound cond.

$$t = \langle w_i | H_1 | w_j \rangle$$

$$U = \langle w_i^2 | V_{12} | w_i^2 \rangle$$



## Renormalized wave equation:

$$\frac{\delta(E - \mu N_e)}{\delta w_i^*(\mathbf{r})} - \nabla \cdot \frac{\delta(E - \mu N_e)}{\delta(\nabla w_i^*(\mathbf{r}))} = \sum_{i \geq j} \lambda_{ij} w_j(\mathbf{r})$$

Adjustable Slater or STO-3G  
basis forms a trial Wannier  
function obtained variationally

## Atomic functions:

$$\Phi_i(\mathbf{r}) = (\pi \alpha^3)^{1/2} \exp(-\alpha |\mathbf{r} - \mathbf{R}_i|)$$

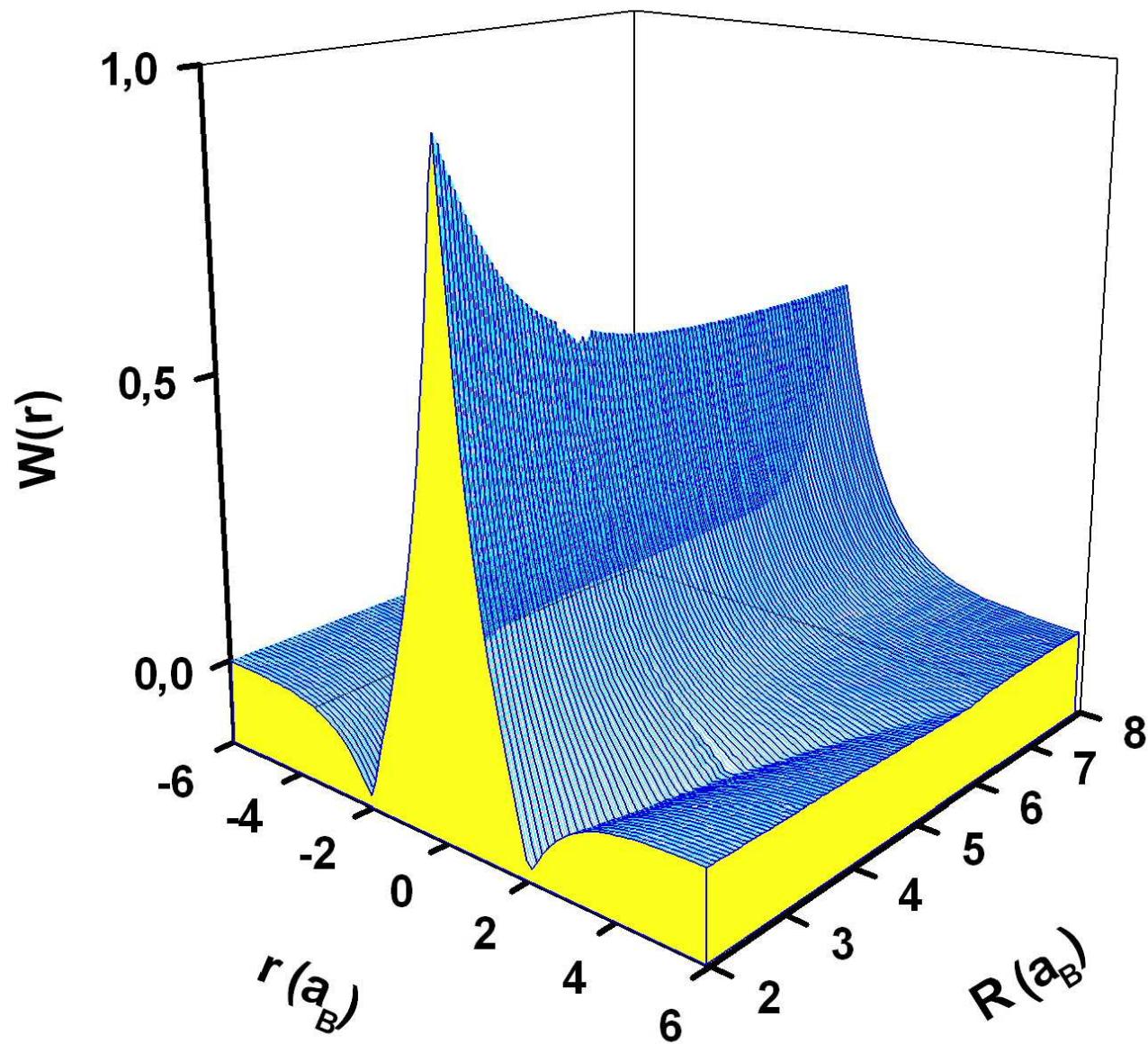
$$\langle \Phi_i | \Phi_j \rangle = S_{ij}$$

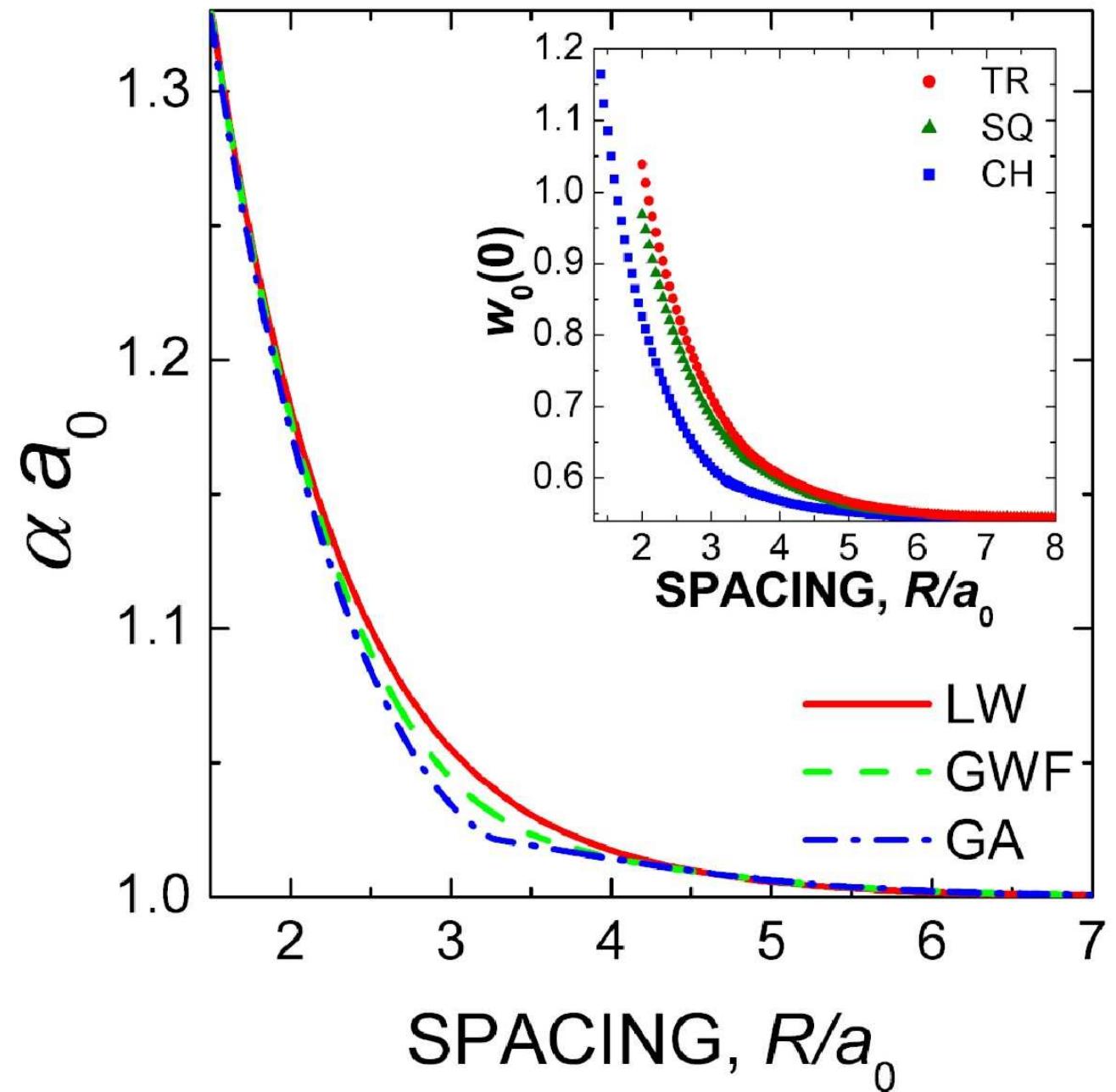
## Wannier functions (wave functions):

$$w_i(\mathbf{r}) = \sum_j \beta_{ij} \Psi_j(\mathbf{r})$$

$$\langle w_i | w_j \rangle = \delta_{ij}$$

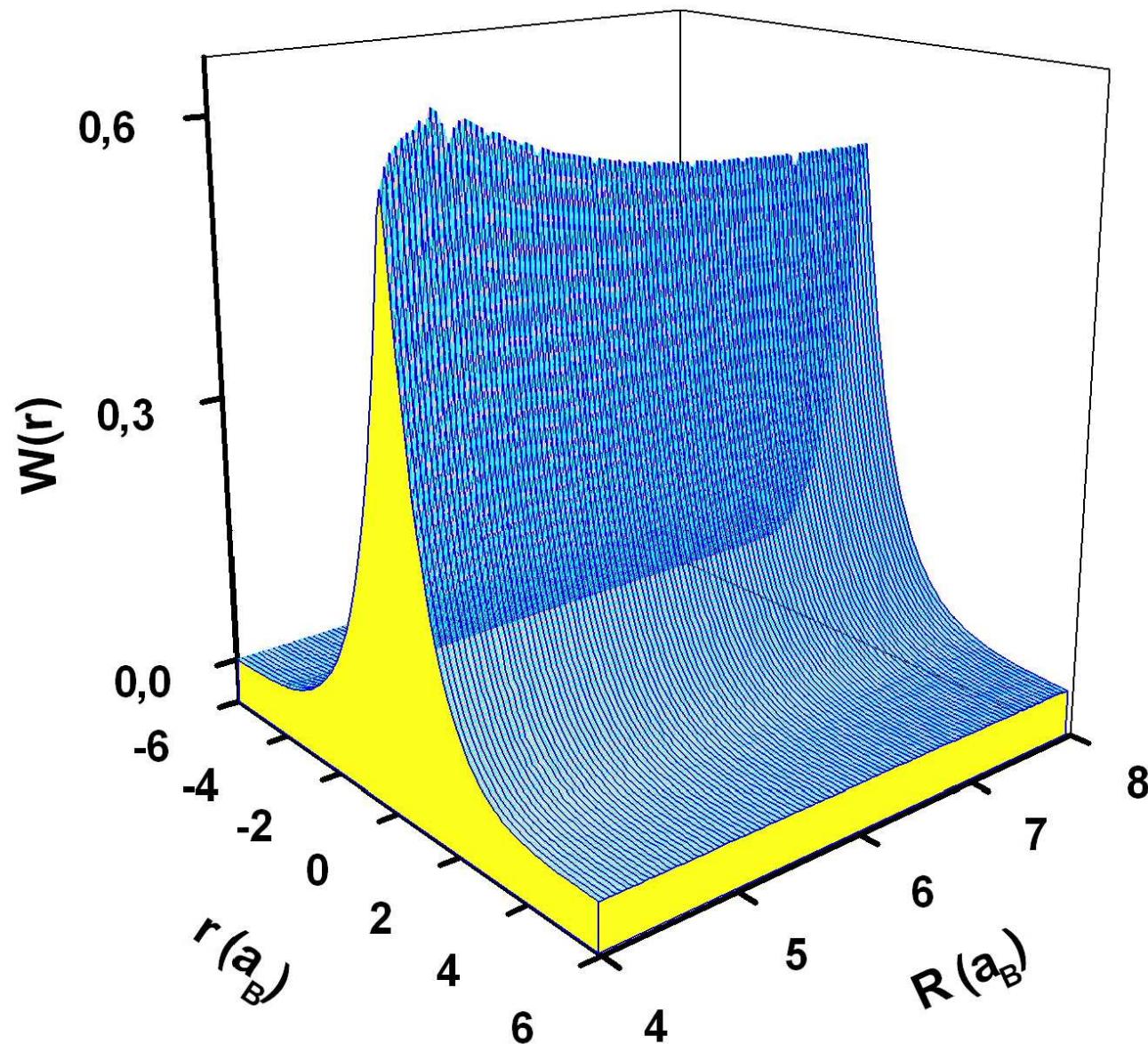
# Square lattice

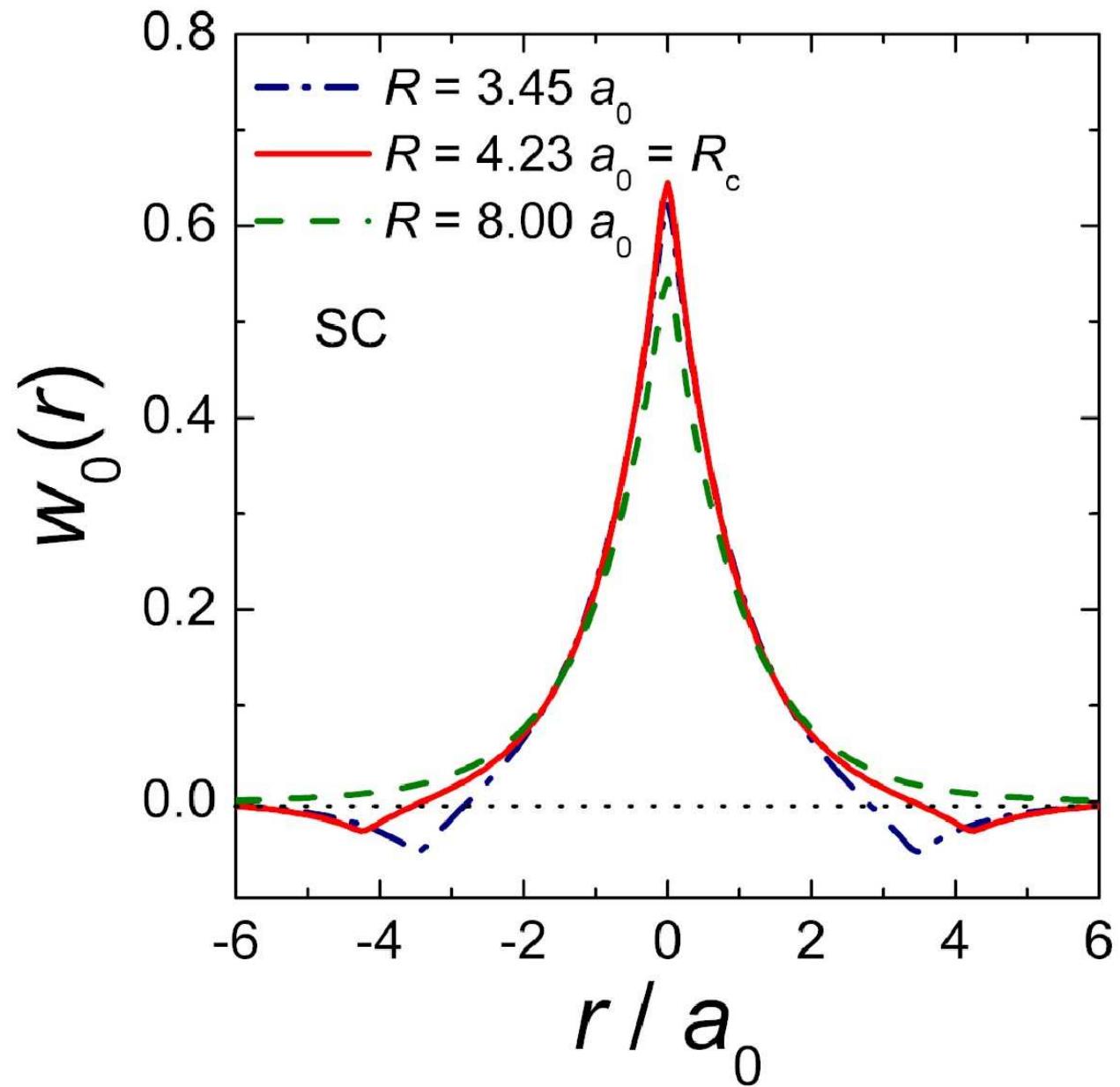


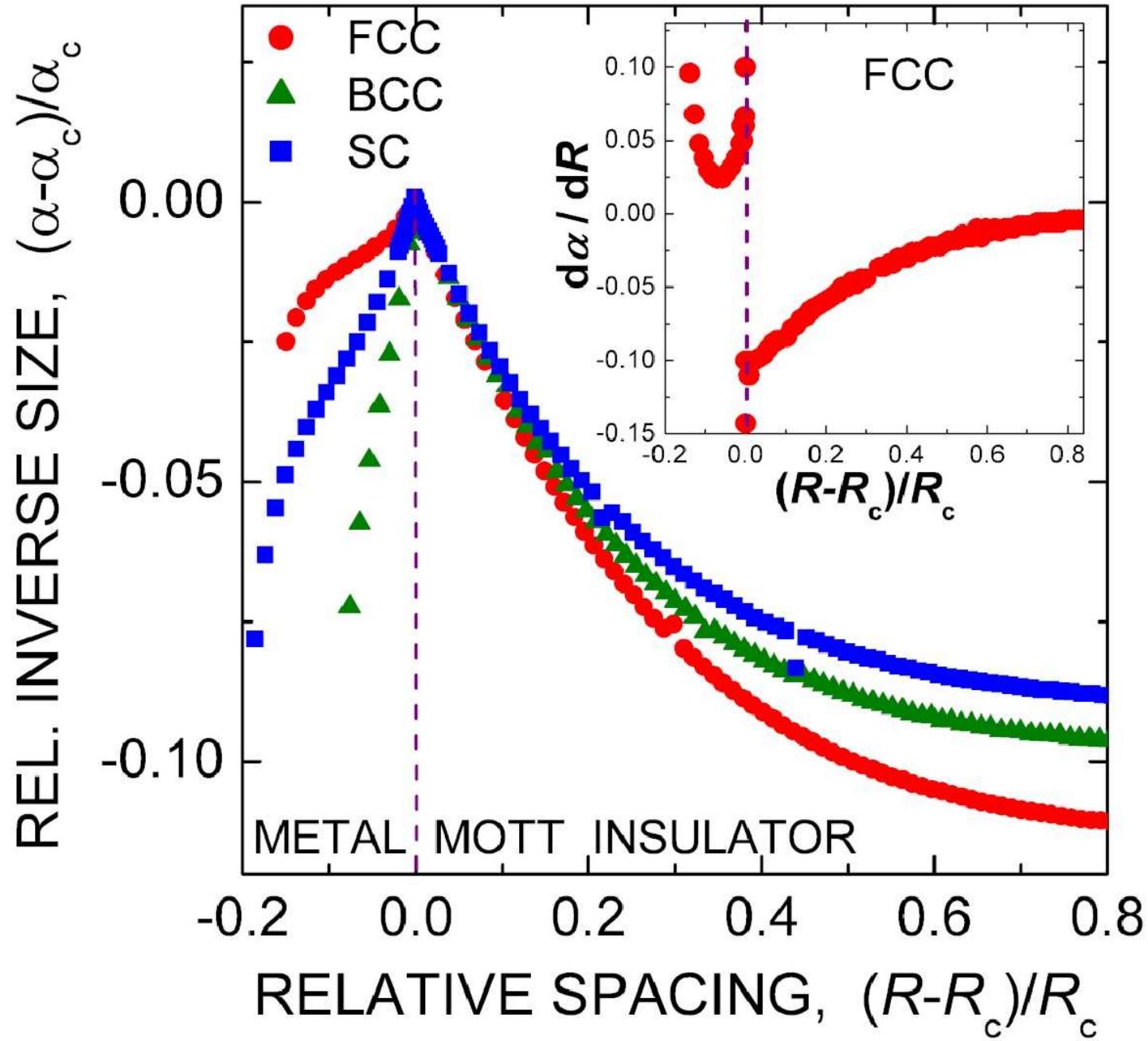


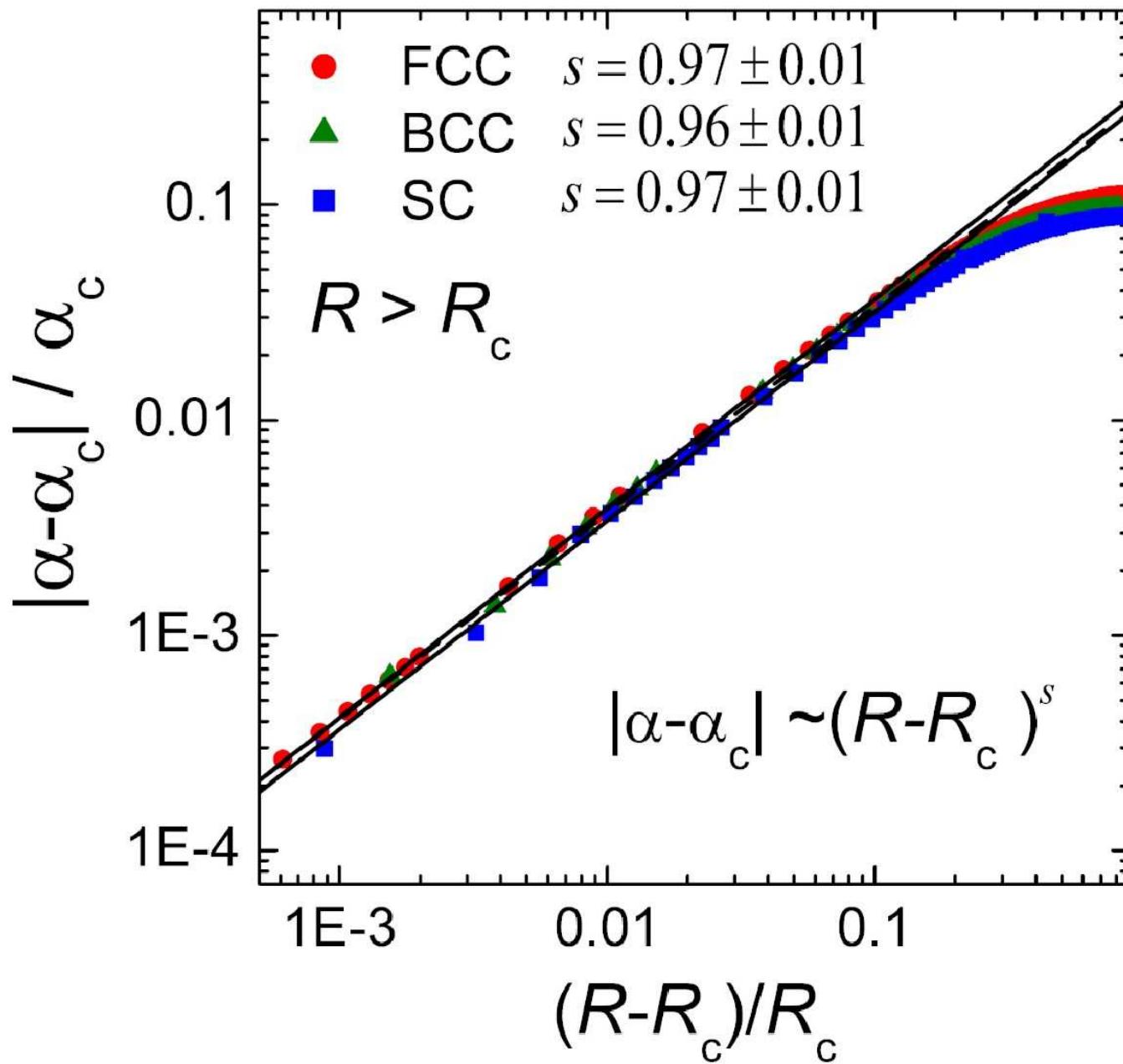
# *3 Dimensions: Gutzwiller approach*

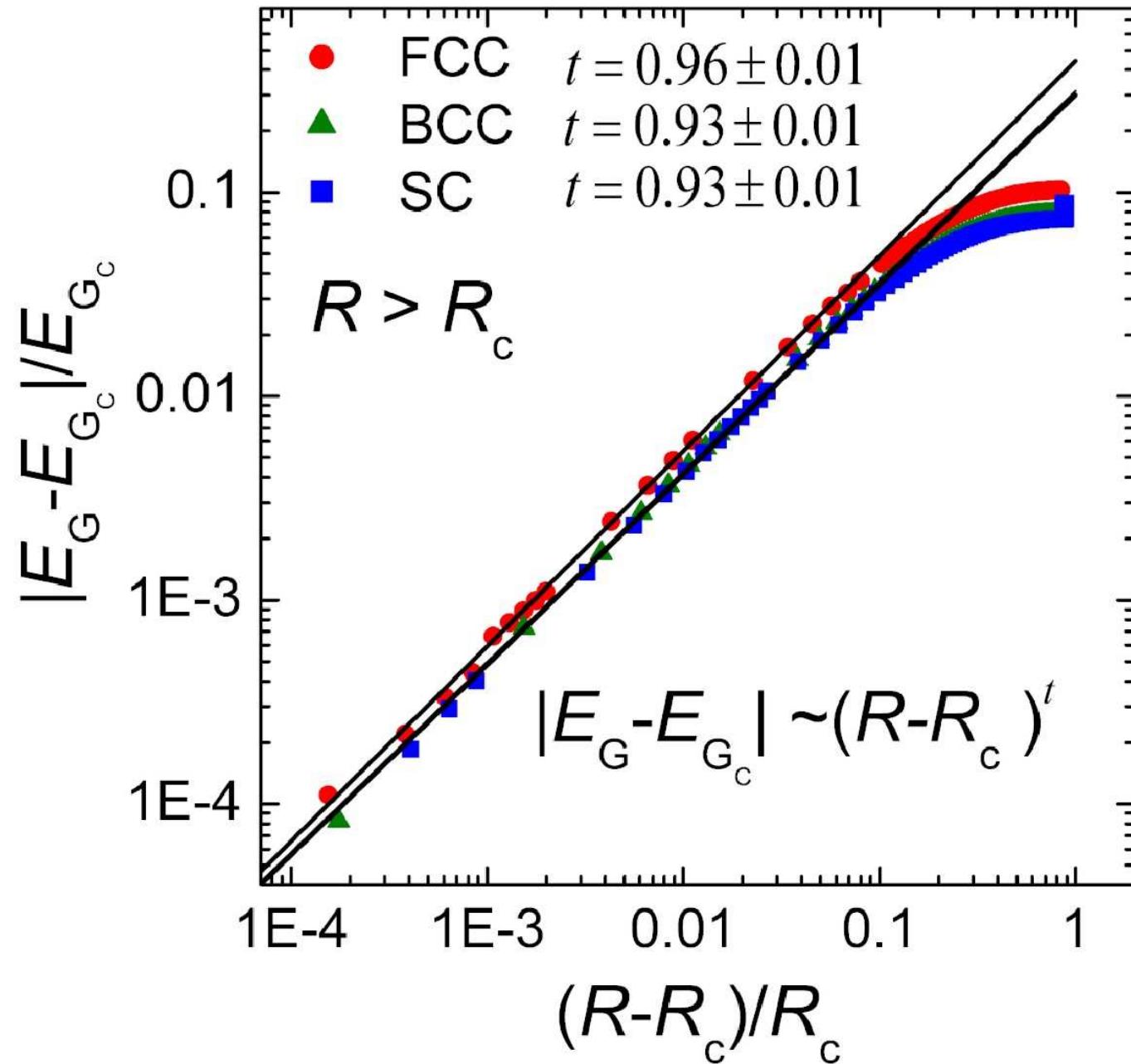
# Body Centered Cubic lattice

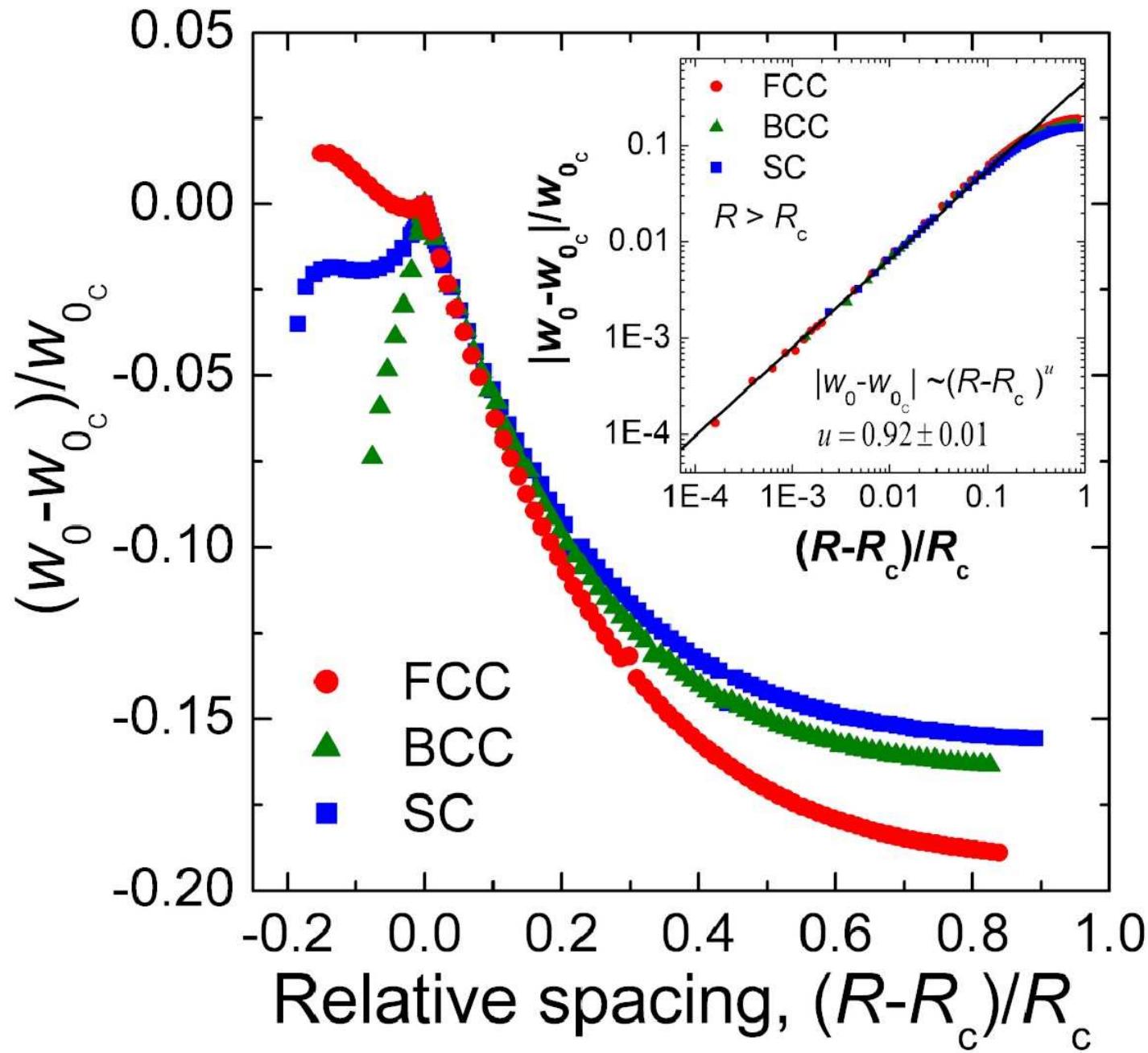












# **Outlook**

- 1. Method allows for study of the electron state evolution as a function of interatomic spacing**
- 2. The evolution of the wave function in the correlated state through the Mott threshold:  
from atoms to solid state (or vice versa)**
- 3. Scaling and critical behavior of the wave function and a critical behavior**
- 4. Future: Bose Hubbard d-orbitals**